

Hybrid modelling based on two-phase flow analogy: advances and challenges

Sergey Karabasov

Ivan Korotkin, Dmitry Nerukh, Elvira Tarasova, Vladimir Farafonov,
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School of Engineering and Materials Science

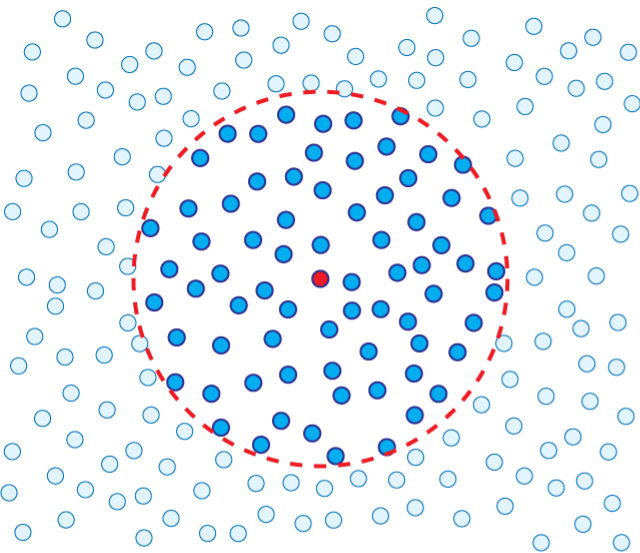
Acknowledgement

- Engineering and Physical Sciences Research Council
- European Commission (MC IF HIPPOGRIFFE project starting July 2016)
- A. Scukins, V. Glotov, V. Goloviznin, V. Semiletov, J. Hu, P.Vadgama

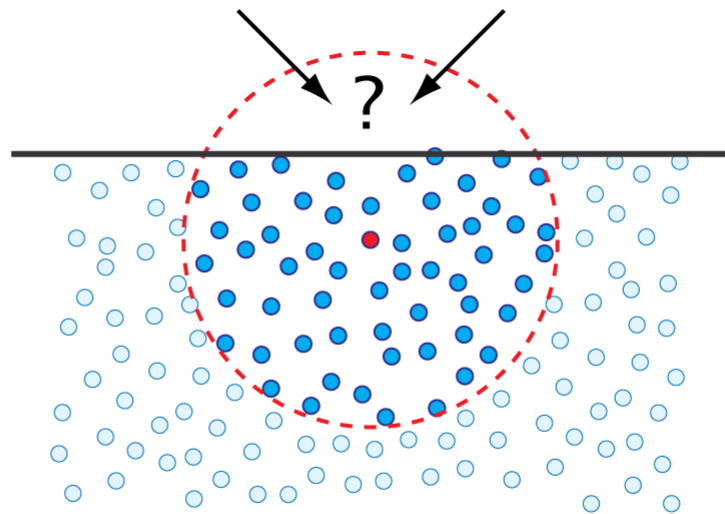
Aston University,
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Continuum \Leftrightarrow Atoms ??

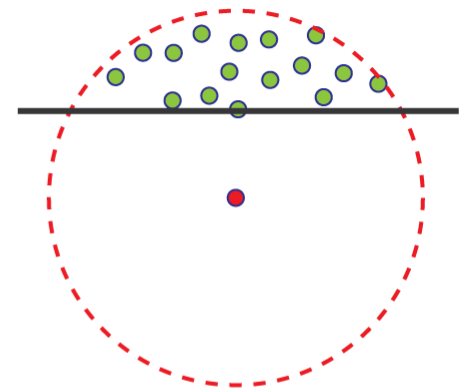
- **Problem:** Missing interactions near “wall”
 - Results in unnatural “wiggles”
- **Possible** solutions to mimic the missing force:
 - Average force normally “felt” by particle
 - Value of force function of distance to wall



Particle in bulk
all force interactions OK

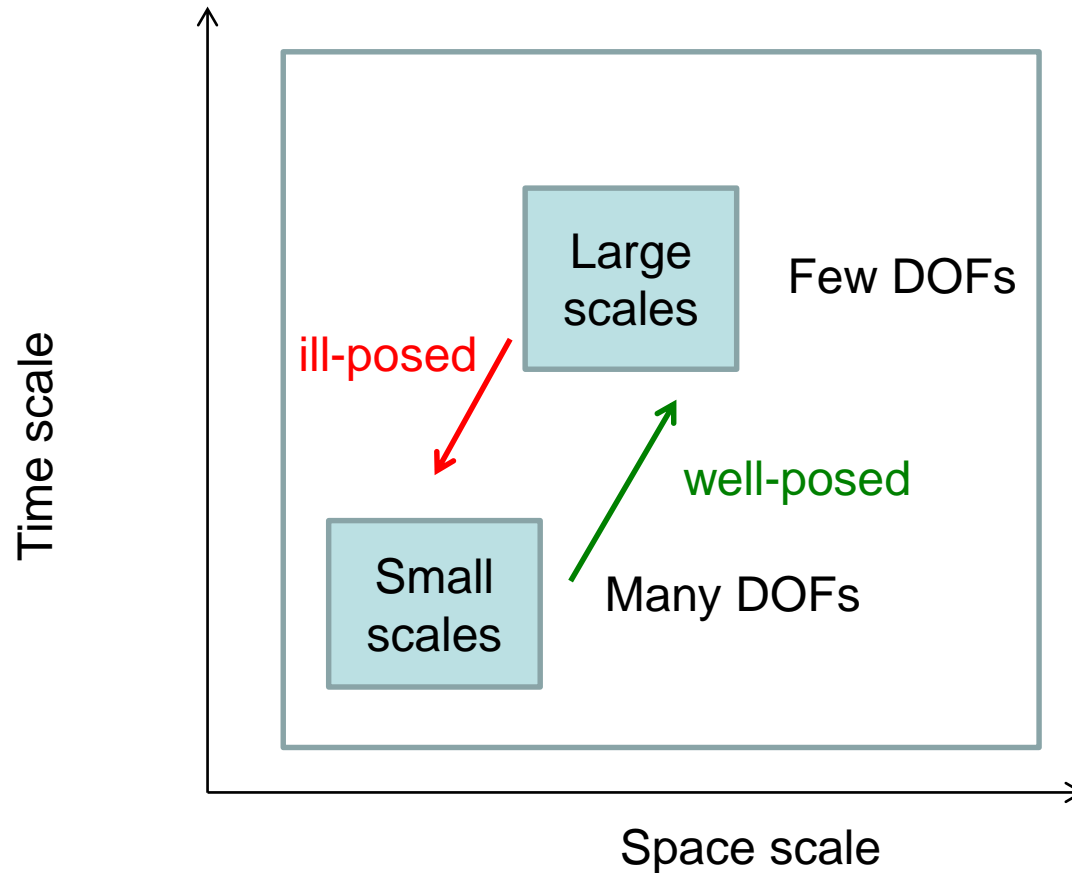


Particle near specular wall
missing force interactions due to wall



Missing part
try to mimic this force

Multiscale modelling: cyclic (fully coupled) approach



Physical analogies as a multiscale modelling method

Examples from physics: electrical circuits, acoustics...

pros: may allow efficient solutions for very complex systems, may be “physically insightful” (depending on the question asked)

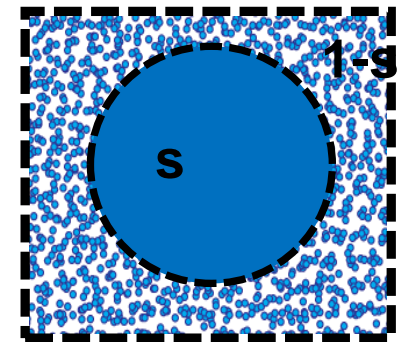
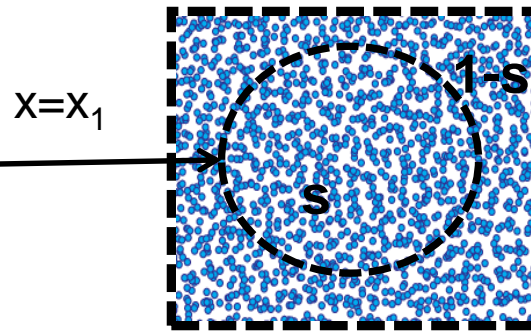
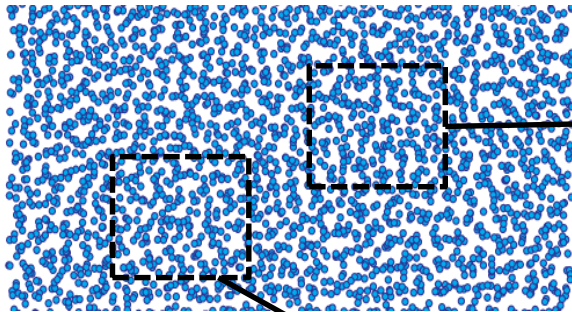
cons: the solution is not unique, not from the “first principles”, “too engineering”, etc..

MD/FH coupling based on the two-phase analogy

S.E. Buckley and M.C. Leverett (1942). "Mechanism of fluid displacements in sands"

Original MD system

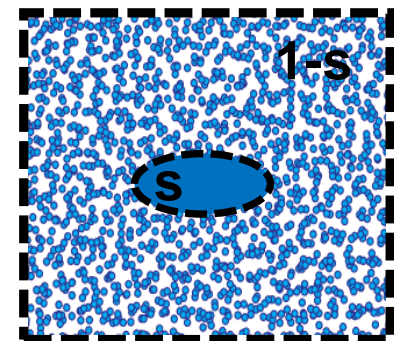
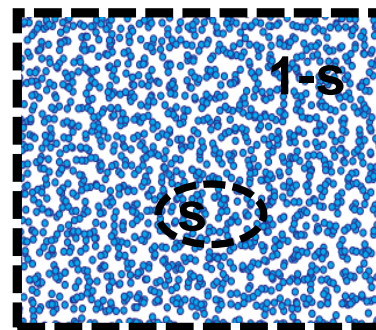
Unit Eulerian volume (x) and partial concentration $S=S(x)$



s =LARGE SCALES

$1-s$ =Small scales

$x=x_2$



$1-S$ – partial volume occupied by MD model

S – partial volume occupied by continuum model

The two 'phases' occupy the same elementary volume of the same liquid, no interface forces are relevant

Two-phase hydrodynamic analogy: mass conservation

S.E. Buckley and M.C. Leverett (1942). "Mechanism of fluid displacements in sands"

Continuum phase

$$\delta_t(sm) + \sum_{\gamma=1,6} (s\rho\bar{\mathbf{u}})d\mathbf{n}^\gamma dt = \delta_t J^{(\rho)}$$

Particle phase

$$\delta_t \left((1-s) \sum_{p=1, N(t)} m_p \right) + \sum_{\gamma=1,6} \left((1-s) \sum_{p=1, N_\gamma(t)} \rho_p \mathbf{u}_p \right) d\mathbf{n}^\gamma dt = -\delta_t J^{(\rho)}$$

Conservation law is satisfied for the mixture density

$$\bar{\rho} = s\rho + (1-s) \sum_{p=1, N(t)} \rho_p$$

J is the model birth/death function that depends on S for the solution to satisfy compatibility conditions: for S-> 1 all phases -> continuum phase, for S-> 0 all phases -> atomistic phase

Two-phase hydrodynamic analogy : momentum conservation

S.E. Buckley and M.C. Leverett (1942). "Mechanism of fluid displacements in sands"

Continuum phase

Landau-Lifshitz'

deterministic + stochastic stresses

$$\delta_t(smu_i) + \sum_{\gamma=1,6} (s\rho u_i \bar{\mathbf{u}}) d\mathbf{n}^\gamma dt = s \sum_{j=1,3} \sum_{\gamma=1,6} (\Pi_{ij} + \tilde{\Pi}_{ij}) dn_j^\gamma dt + \delta_t J_i^{(\mathbf{u})} dt, i = 1,3$$

Particle phase

$$\delta_t \left((1-s) \sum_{p=1, N(t)} m_p u_{ip} \right) + \sum_{\lambda=1,6} \left((1-s) \sum_{p=1, N_\lambda(t)} \rho_p u_{ip} \mathbf{u}_p \right) d\mathbf{n}^\lambda dt = (1-s) \sum_{p=1, N(t)} F_{ip} dt - \delta_t J_i^{(\mathbf{u})} dt, i = 1,3$$

2nd Newton's law is satisfied: the change of the mixture momentum

$$\bar{\rho} \cdot \bar{u}_i, \text{ where } \bar{u}_i = \left[s\rho u_i + (1-s) \sum_{p=1, N(t)} \rho_p u_{ip} \right] = \text{force}$$

Closing the hybrid model

Specify the source birth/death terms in the 'buffer zone' $0 < s < 1$
the continuum solution is **forced (softly diffused or exponentially forced)** towards the 'target' MD solution;

$$D_t \left(\bar{m} - \sum_{p=1, N(t)} m_p \right) = L^{(\rho)} \bullet \left(\bar{m} - \sum_{p=1, N(t)} m_p \right),$$
$$D_t \left(\bar{u}_i \bar{m} - \sum_{p=1, N(t)} u_{ip} m_p \right) = L^{(u)} \bullet \left(\bar{u}_i \bar{m} - \sum_{p=1, N(t)} u_{ip} m_p \right) + s \sum_{j=1,3} \sum_{\gamma=1,6} (\Pi_{ij} + \tilde{\Pi}_{ij}) dn_j^\gamma dt,$$

for $s=1$ and $s=0$ allow the model to become pure continuum (fluctuating hydrodynamics) and pure MD, respectively.

Consistent modification of the MD equations

- Add forcing terms to the molecular dynamics kinematic and dynamic equations to satisfy the macroscopic equations of the hybrid model

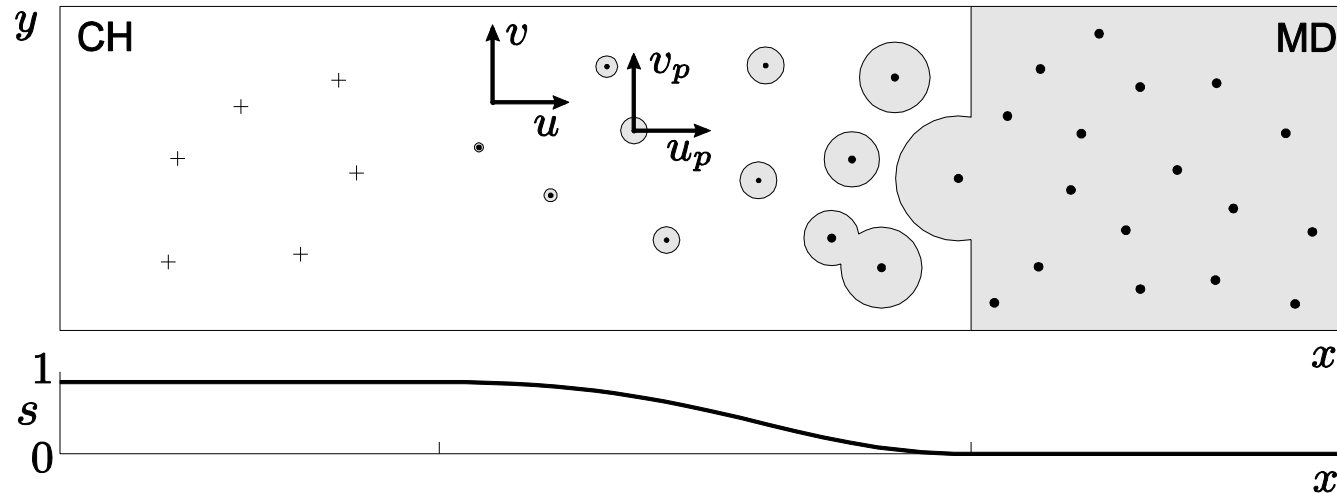
$$\frac{dx_{ip}}{dt}^{MD} = u_{ip}^{Newton} + ?..$$

$$\frac{d}{dt} u_{ip}^{Newton} = - \frac{d}{dx_i} V_p^{MD} + ?..$$

Can work out the expressions for the “?” terms from the corresponding discrete conservation laws for MD particles

Note that $dx_p/dt \neq u_p$ for the “two phase flow” model

2D examples of two-way coupled hybrid models



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OF

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SOCIETY

A

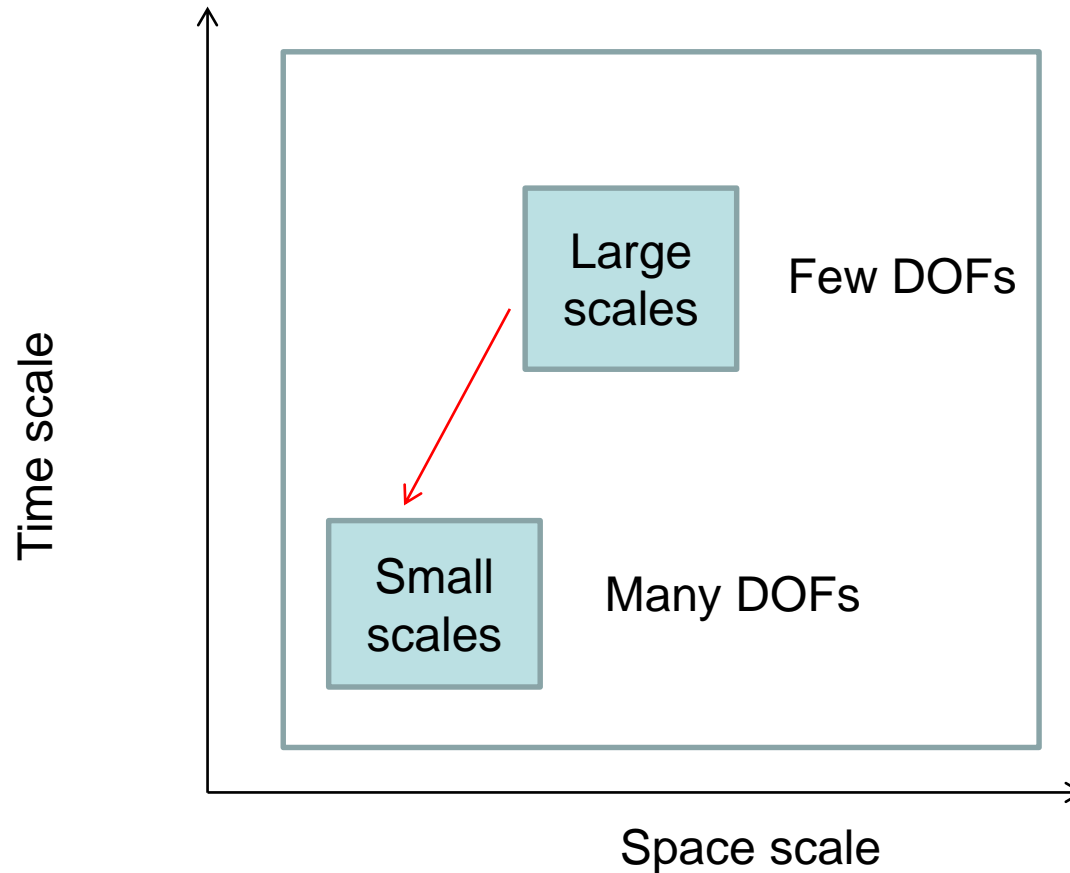
MATHEMATICAL,
PHYSICAL
& ENGINEERING
SCIENCES

Concurrent multiscale modelling of atomistic and hydrodynamic processes in liquids

Anton Markesteijn, Sergey Karabasov, Arturs Scukins, Dmitry Nerukh, Vyacheslav Glotov and Vasily Goloviznin

Phil. Trans. R. Soc. A 2014 **372**, 20130379, published 30 June 2014

(Current) Multiscale modelling in 3D: acyclic “top-bottom”/ “boundary-condition” approach



Effectively works like an ‘open-domain’ boundary condition in CFD

Large scales: Fluctuating Hydrodynamics

- Fluctuating Hydrodynamics
 - Dissipative fluxes treated as stochastic variables
 - Random variables, mimicking molecular motion
- Fluctuation-Dissipation theorem
- Equations of Fluctuating Hydrodynamics
 - Conservation of Mass / Conservation of Momentum
 - Added fluctuating stress tensor

$$\langle \delta \Pi_{ij}(\mathbf{r}, t) \cdot \delta \Pi_{kl}(\mathbf{r}', t') \rangle = 2k_B T \left[\eta (\delta_{ik} \delta_{jl} + \delta_{il} \delta_{jk}) + \left(\eta_V - \frac{2}{3} \eta \right) \delta_{ij} \delta_{kl} \right] \times \delta(\mathbf{r} - \mathbf{r}') \delta(t - t').$$

Landau-Lifshitz Fluctuating Hydrodynamics Equations

One dimensional case

$$\frac{\partial \rho}{\partial t} + \frac{\partial \rho u}{\partial x} = 0$$

$$\frac{\partial \rho u}{\partial t} + \frac{\partial (\rho u^2 + P)}{\partial x} - \frac{4}{3} \cdot \eta \cdot \frac{\partial^2 u}{\partial x^2} - \frac{\partial s}{\partial x} = 0$$

$$\frac{\partial \rho E}{\partial t} + \frac{\partial (\rho E + P)u}{\partial x} - \frac{\partial}{\partial x} \left(\frac{4}{3} \cdot \eta u \frac{\partial u}{\partial x} + \kappa \cdot \frac{\partial T}{\partial x} \right) - \frac{\partial (q + u \cdot s)}{\partial x} = 0$$

Stochastic fluxes

$$\langle s(x,t) s(x',t') \rangle = \frac{8 \cdot k \cdot \eta \cdot T}{3 \cdot \sigma} \cdot \delta(x-x') \cdot \delta(t-t');$$

$$\langle q(x,t) q(x',t') \rangle = \frac{2 \cdot k \cdot \kappa \cdot T^2}{\sigma} \cdot \delta(x-x') \cdot \delta(t-t');$$

$$\rho E = c_v \rho T + \frac{\rho u^2}{2};$$



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Comput. Methods Appl. Mech. Engrg. 281 (2014) 29–53

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A new non-linear two-time-level Central Leapfrog scheme in staggered conservation–flux variables for fluctuating hydrodynamics equations with GPU implementation

A.P. Markesteijn^{a,*}, S.A. Karabasov^a, V.Yu. Glotov^b, V.M. Goloviznin^b

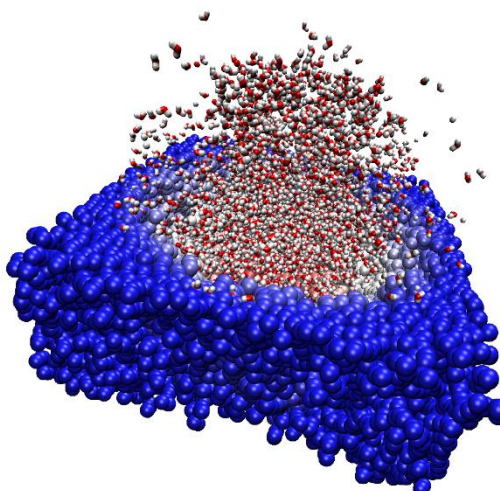
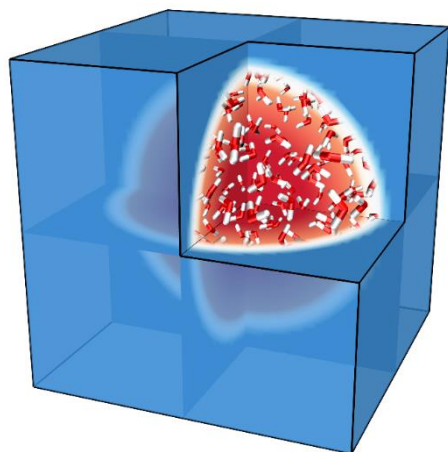
From fluctuating hydrodynamics to all-atom resolution

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A hybrid molecular dynamics/fluctuating hydrodynamics method for modelling liquids at multiple scales in space and time  CrossMark


Ivan Korotkin^{1,a)}, Sergey Karabasov¹, Dmitry Nerukh², Anton Markesteijn¹, Arturs Scukins², Vladimir Farafonov³ and Evgen Pavlov^{2,4}

[+ VIEW AFFILIATIONS](#)

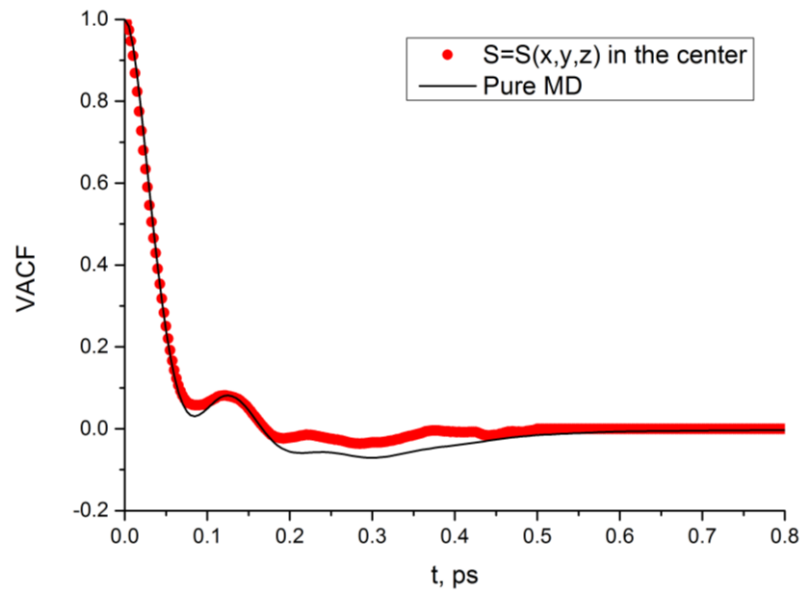
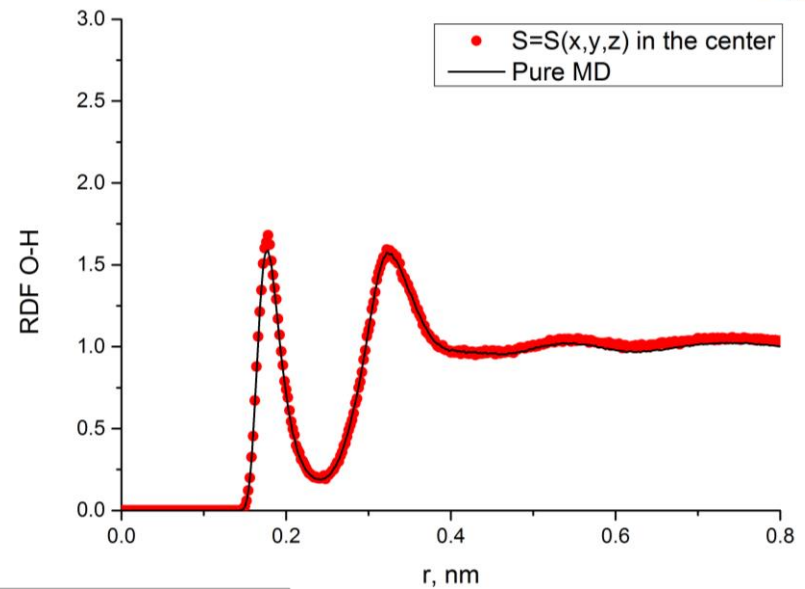
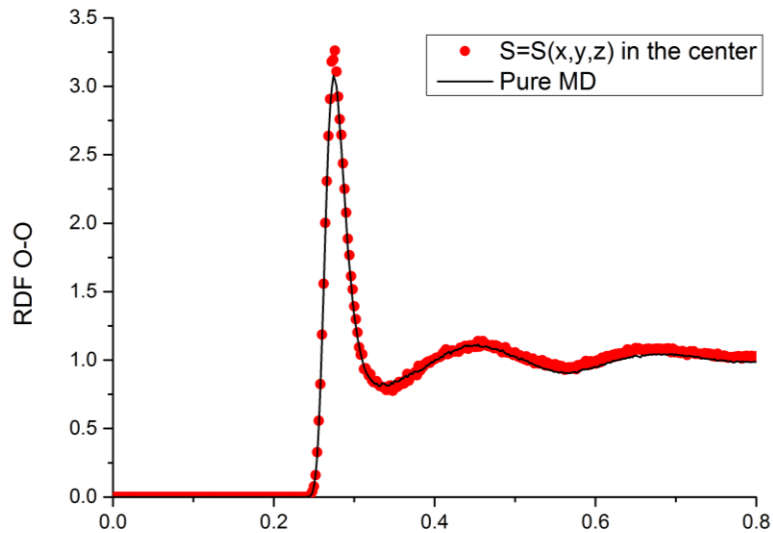
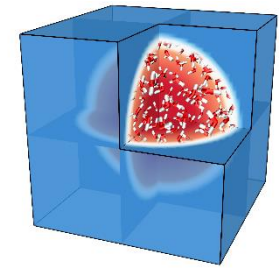
^{a)} Author to whom correspondence should be addressed. Electronic mail: i.korotkin@qmul.ac.uk

J. Chem. Phys. **143**, 014110 (2015); <http://dx.doi.org/10.1063/1.4923011> 

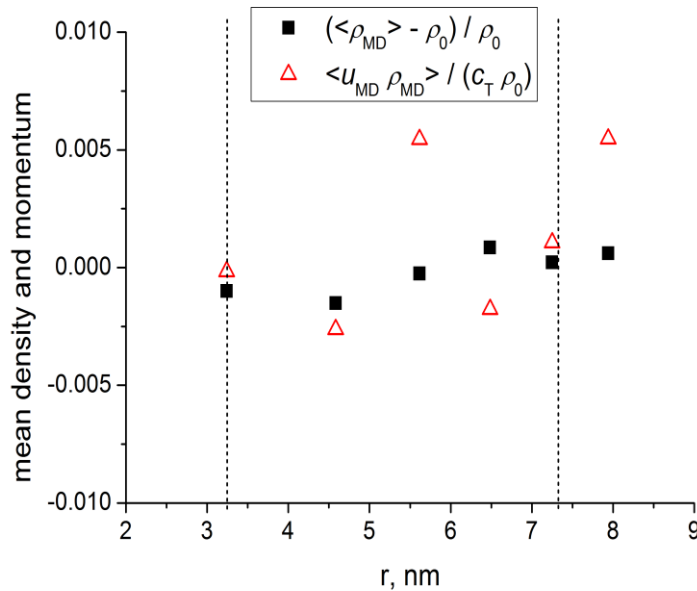
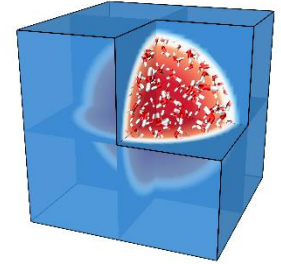
 Buy: \$30.00

 Rent: \$4.00 

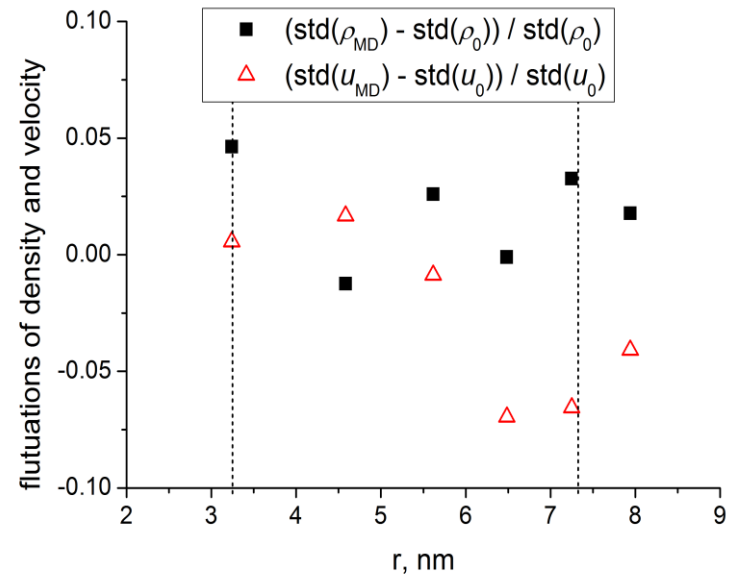
'Simple water': radial distributions and velocity correlations



Preservation of mass, momentum, and correct fluctuation amplitudes across the hybrid zone

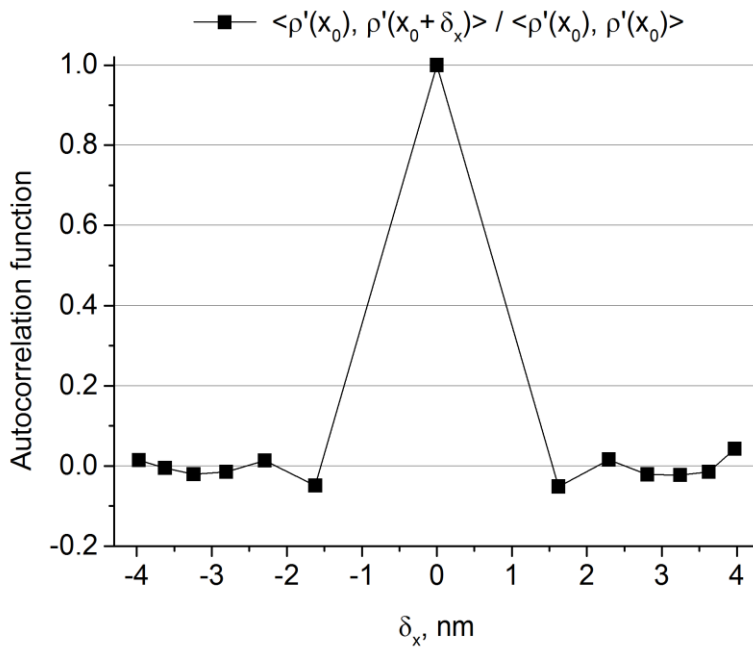
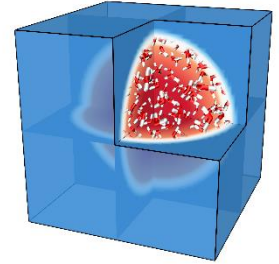


means

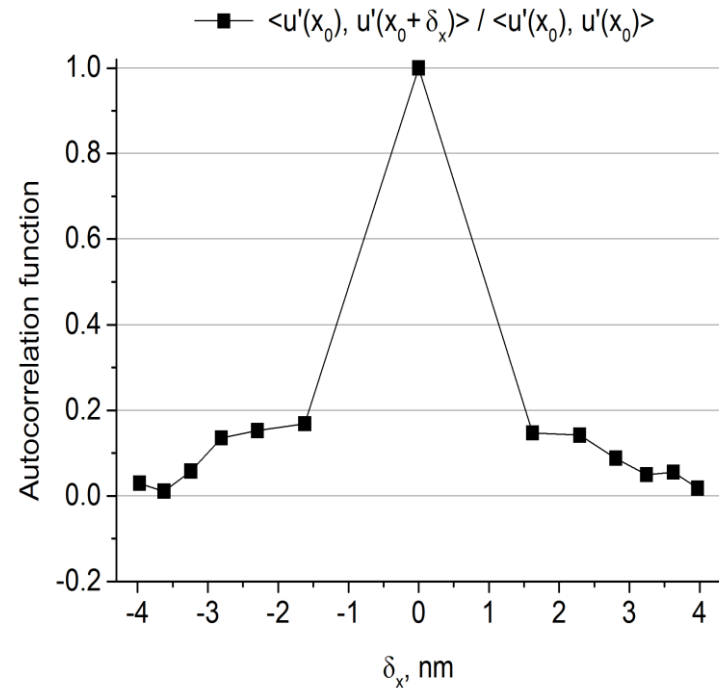


fluctuations

Preservation of random fluctuations across the hybrid zone

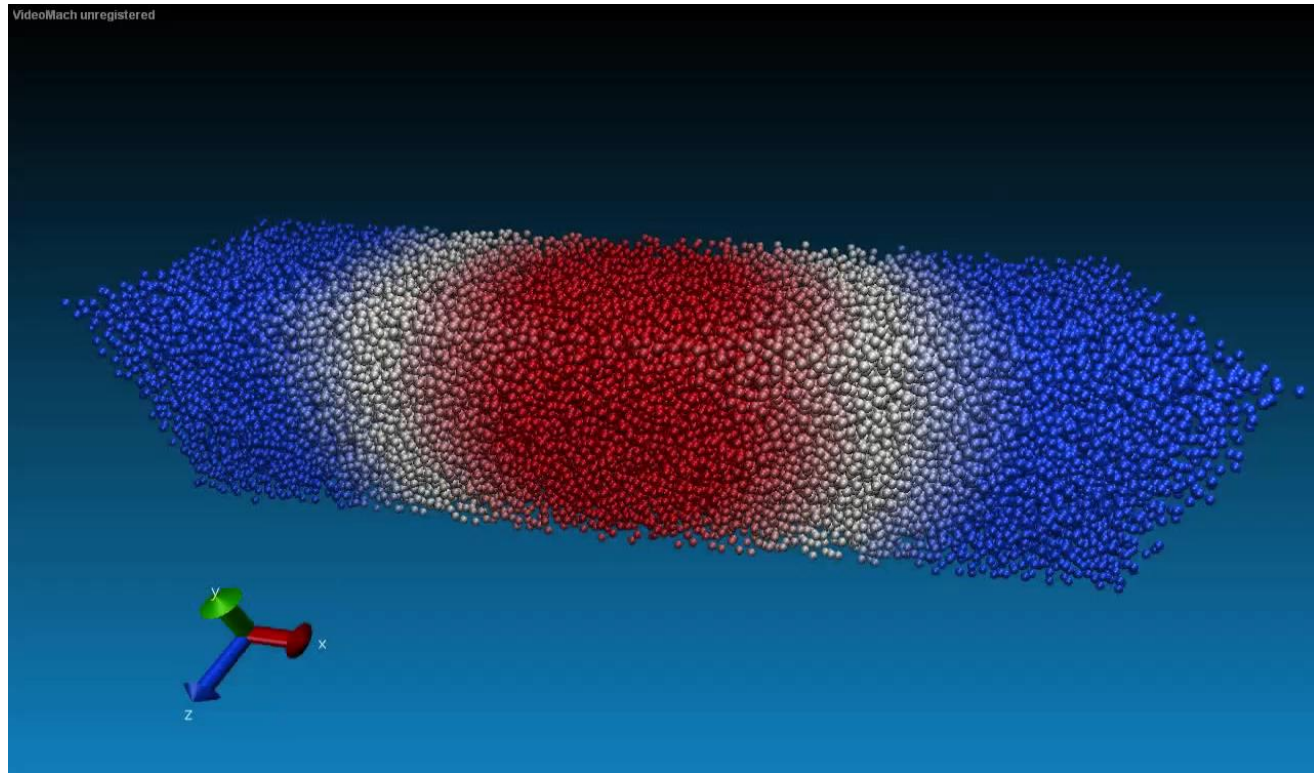


density

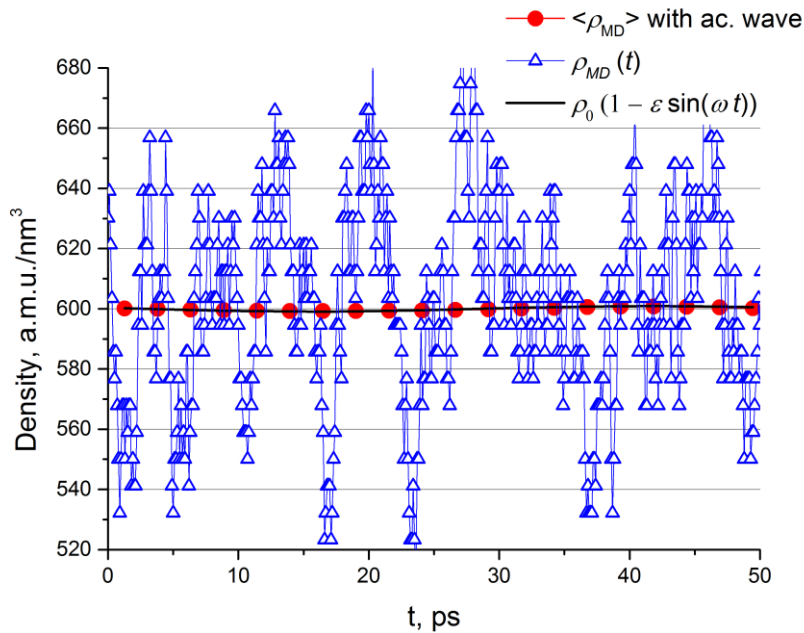
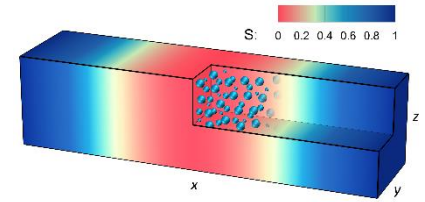


velocity

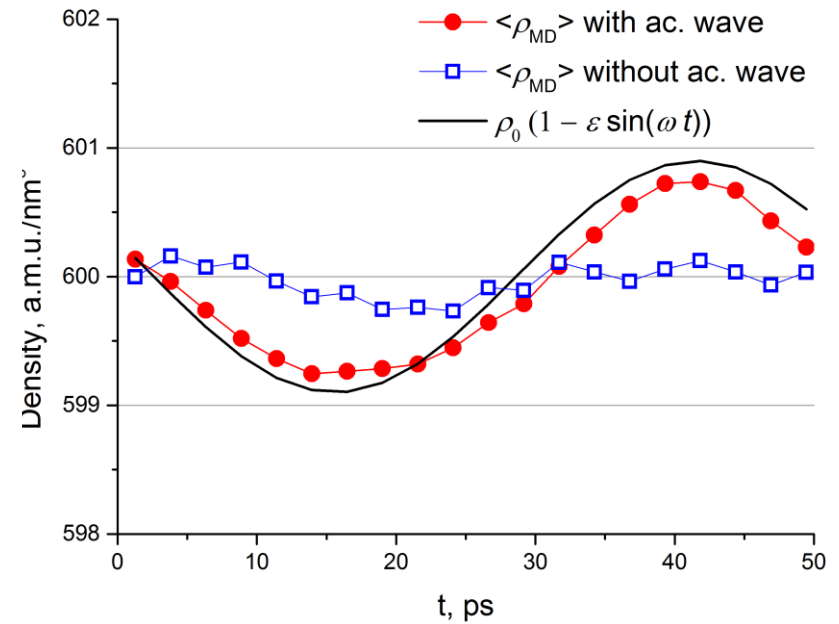
Test problem: hydrodynamic waves travelling through the hybrid MD/FH zone (20x5x5 FH cells)



Case 1: acoustic wave



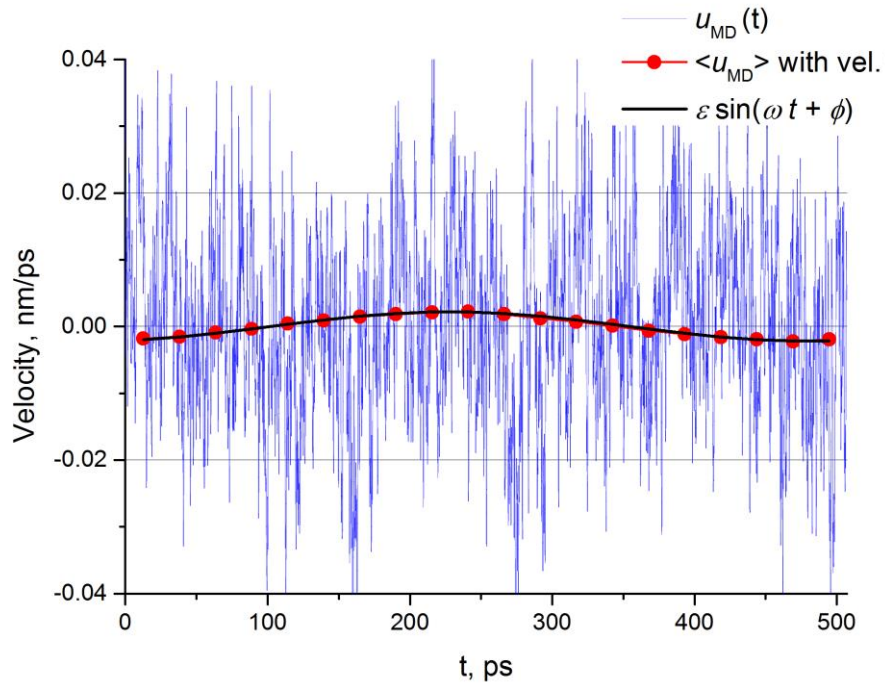
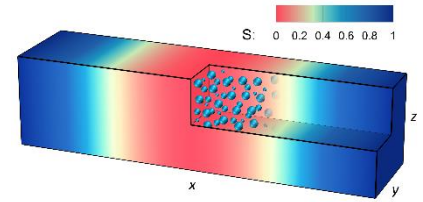
Original signal



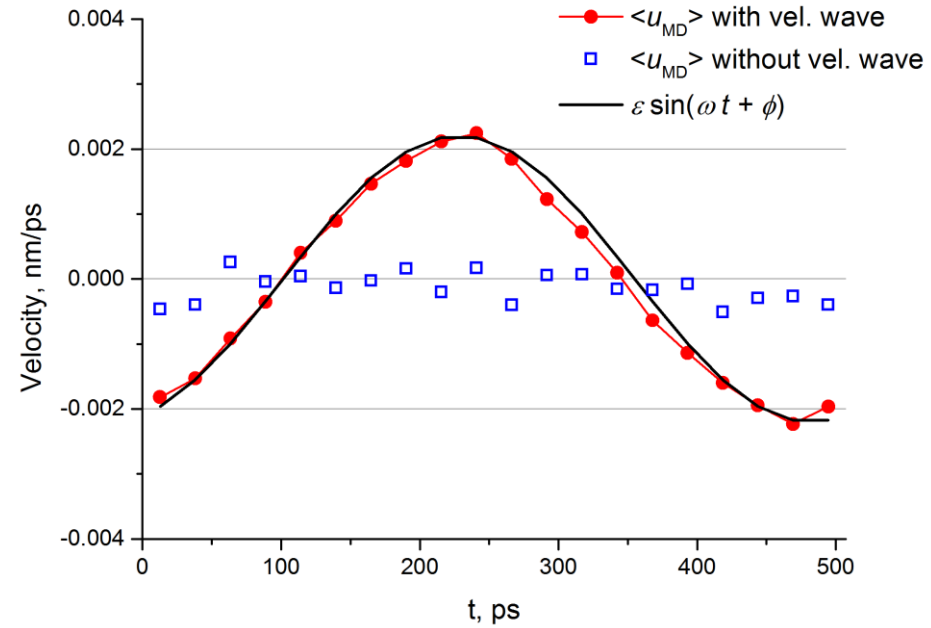
Phase and y&z averaged signal

Signal/noise ~ 0.01

Case 2: shear wave



Original signal



Phase and y&z averaged signal

Signal/noise ~ 0.05

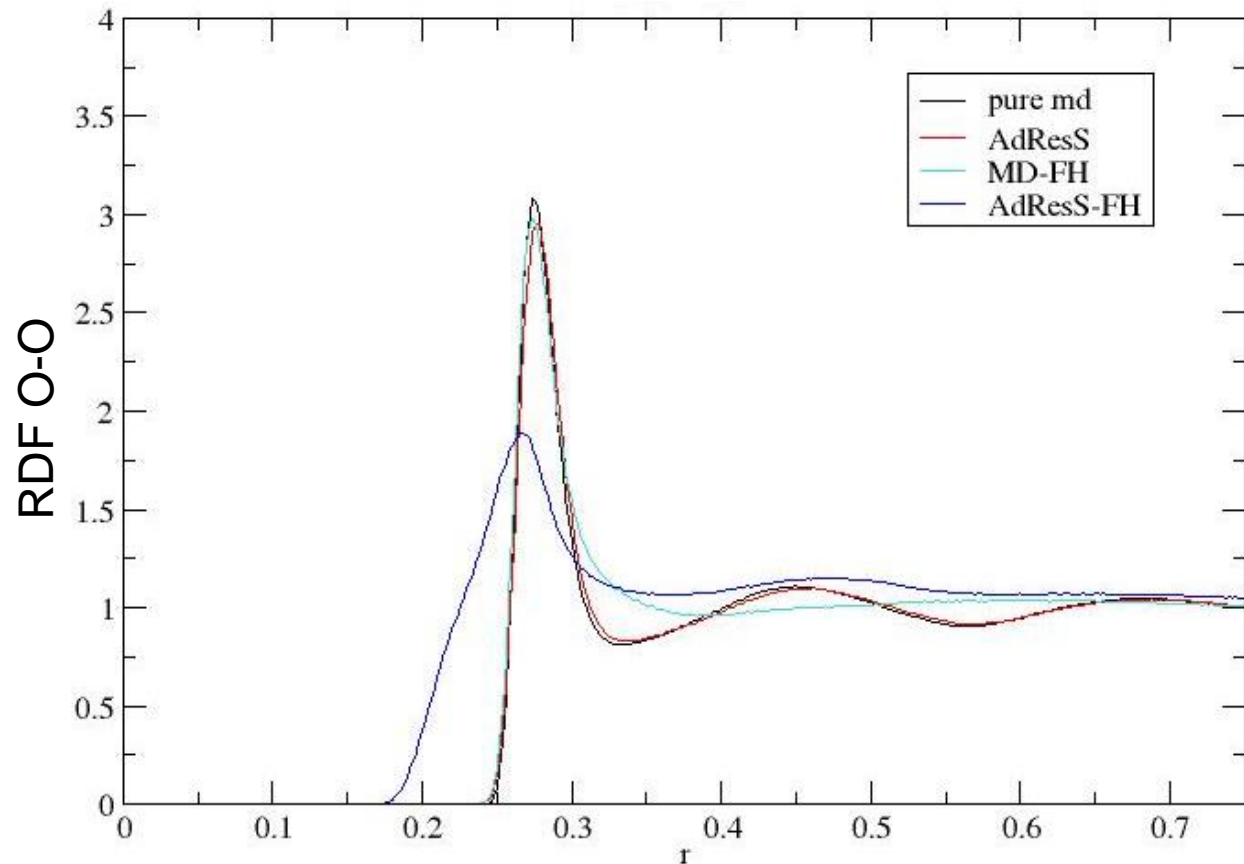
Unsolved issue: sensitivity to the hybrid model parameters

Keep away the redundant MD particles from the FH zone using the space-time expansion ?

Introduce MD (+CG) +FH?



Effect of the macroscopic equation of state?

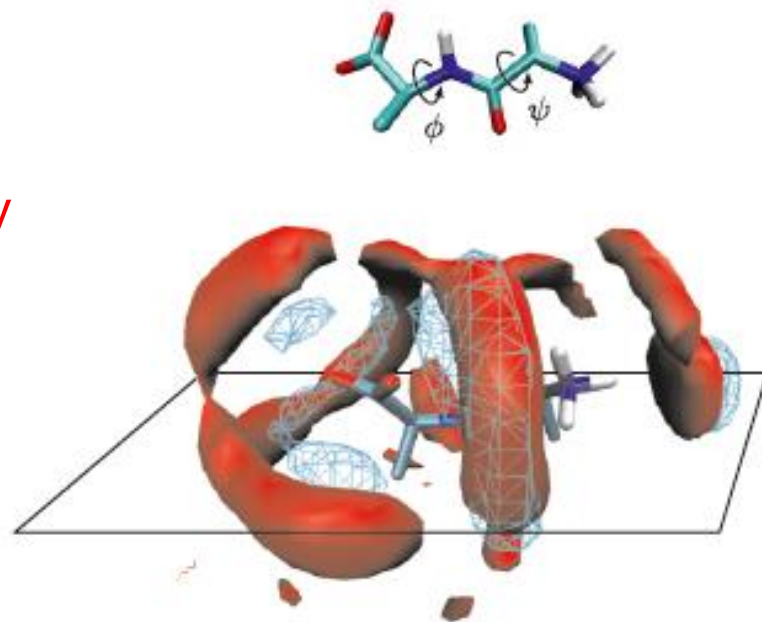


Application: dialanine molecule in water

Water–Peptide Dynamics during Conformational Transitions

Dmitry Nerukh^{*†} and Sergey Karabasov[‡]

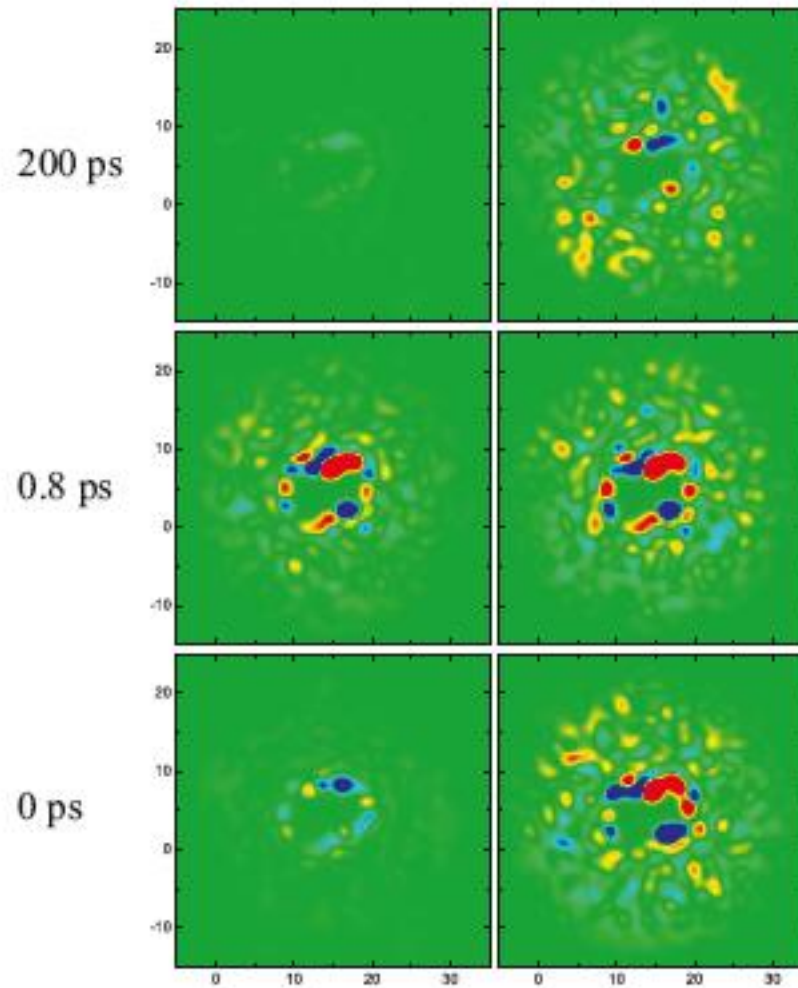
Surrounding water density changes as the dialanine molecule changes its conformation



Question: to what extent the dynamics of water density (fluctuations) is connected with the dynamics of the peptide (quantified by its dihedral angles ϕ and ψ)?

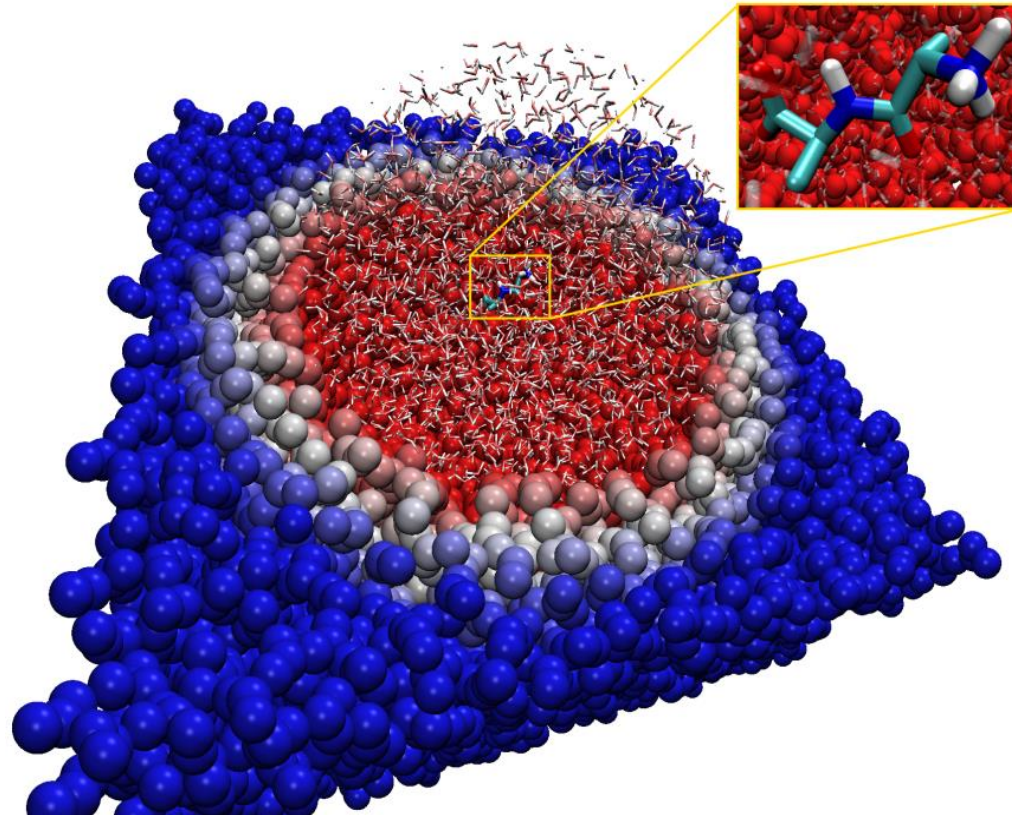
Water density
fluctuations correlated
with the peptide motion

Total
fluctuations



Answer: strongly correlated but *only* at very specific periods, when the conformational transitions occur

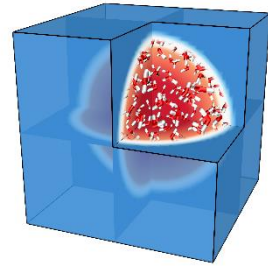
Dialanine in water: a multiscale model



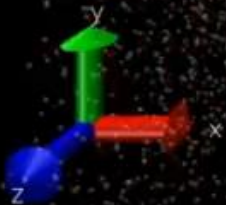
“Static sphere” MD/FH algorithm

- Define S -function, which specifies the (stationary) space decomposition between the MD and the hydrodynamics driven parts of the computational domain,
- For each particle location, find the corresponding control volume it belongs to, interpolate the hydrodynamic fields obtained, compute S -function value at the particle location, and integrate MD equations over one time step,
- Update positions of all particles and repeat.

Dialanine diffusion in water, $s=s(\mathbf{x})$



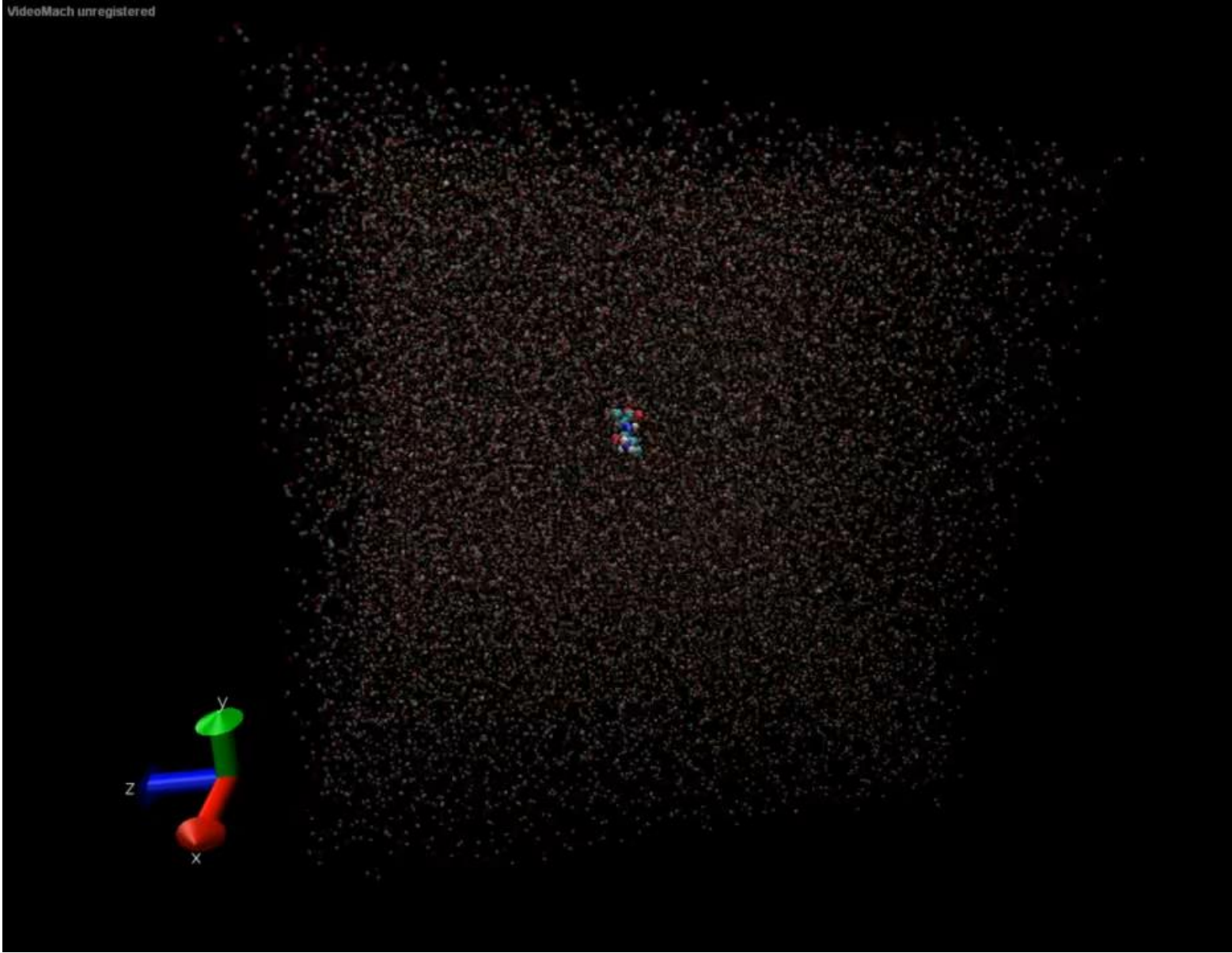
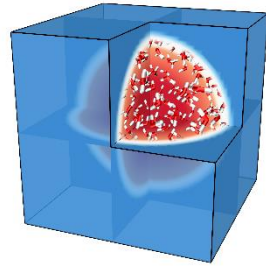
VideoMach unregistered



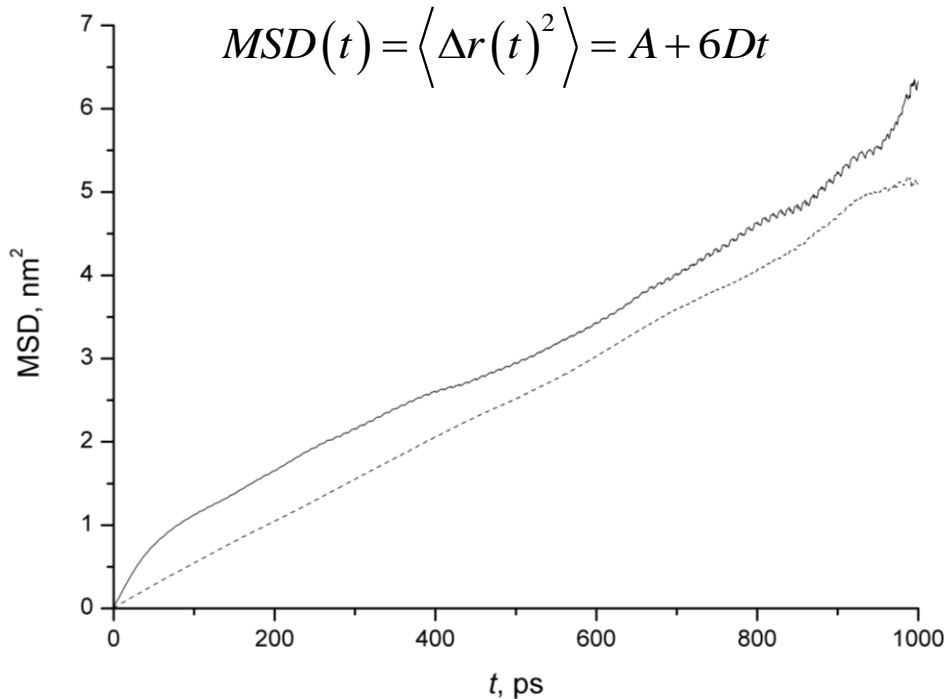
“Dynamic sphere” MD/FH algorithm

- Specify the feature of interest, e.g. the centre of mass of a complex molecule that needs to be treated at all-atom resolution.
- Lock the centre of S-function to this feature so that the coordinates of the origin are updated at each time step.
- For each particle location corresponding to a certain time moment, find the control volume it belongs to, interpolate the hydrodynamic fields to its location, calculate the space-time variable S-function at that location, and integrate MD equations over one time step, then update positions of all particles, move the origin of S-function in accordance with the kinematics of the feature of interest, and repeat.

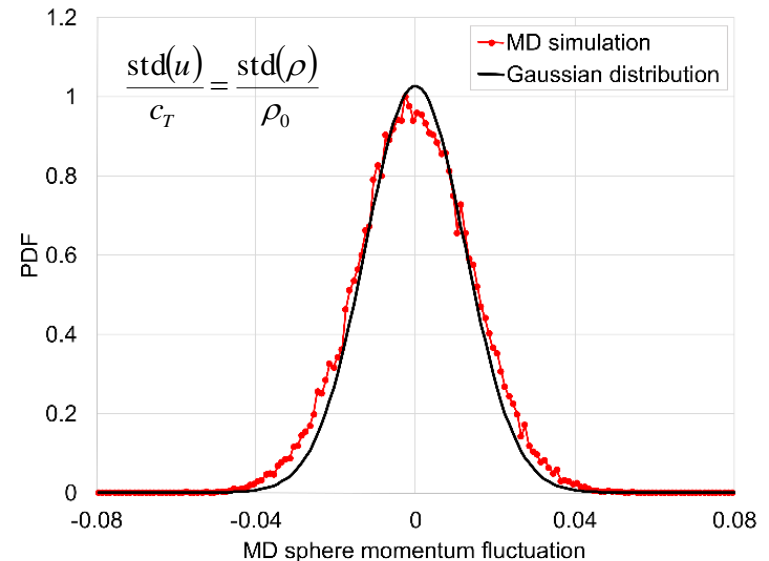
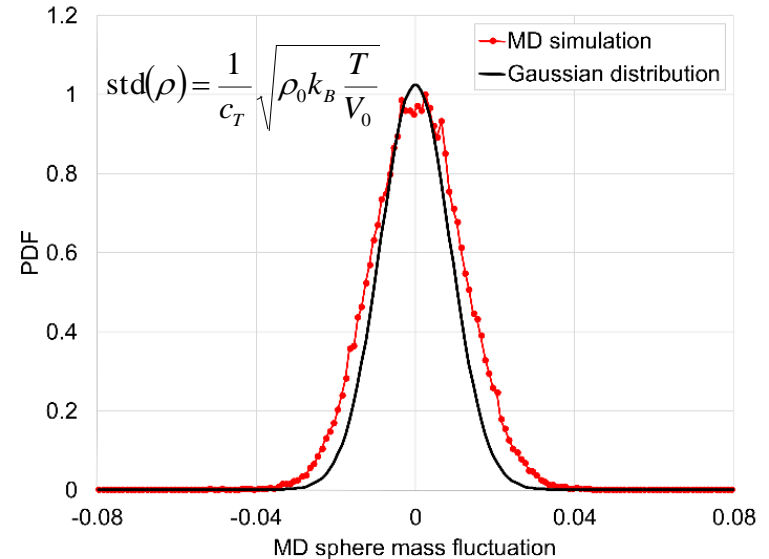
Dialanine diffusion in water, $s=s(\mathbf{x},t)$



Dialanine in water: mass and momentum fluctuations and diffusion coefficient



$$D = (\mathbf{0.83} \pm 0.08) \cdot 10^{-5} \text{ cm}^2/\text{s}$$
$$(D_{\text{ref}} = \mathbf{0.86} \cdot 10^{-5} \text{ cm}^2/\text{s})$$



Application for “Micro/ Nano Fluidics for Health and Diagnostics”?

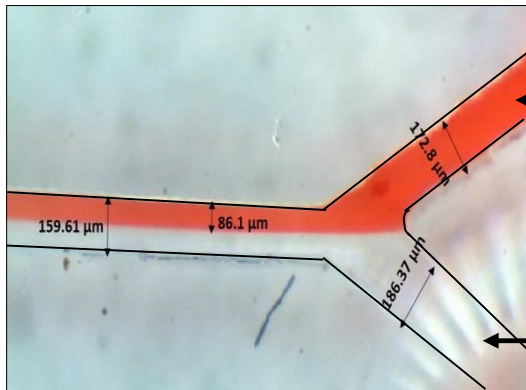
Dr. Ajay Agarwal

- *Principal Scientist, CSIR-CEERI, Pilani, INDIA*
- *Associate Professor, AcSIR, New Delhi, INDIA*

Prof. Pankaj Vadgama

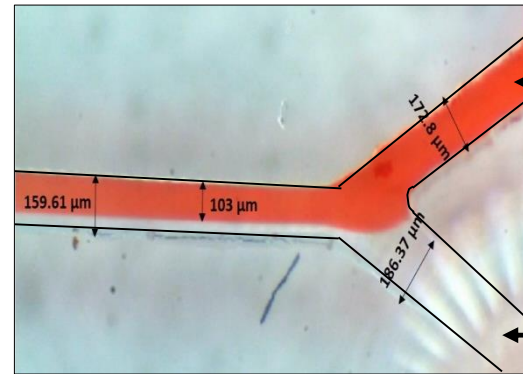
- *Director, IRC in Biomedical Materials,*
 - *Queen Mary Univ. of London, UK*
- **Versatile microfluidic platforms for medical diagnostics**
- **Controlled flow for kinetically and diffusion controlled assays**
- **Dual flow regimens for liquid-liquid interfacing of sample and reagent streams**
- **Integration of sensor arrays into micro flow devices for practical assay and reaction modeling.**
- **Development of pump-less fluid flow using porous absorbent phases**

Effect of flow rates on laminar flow width



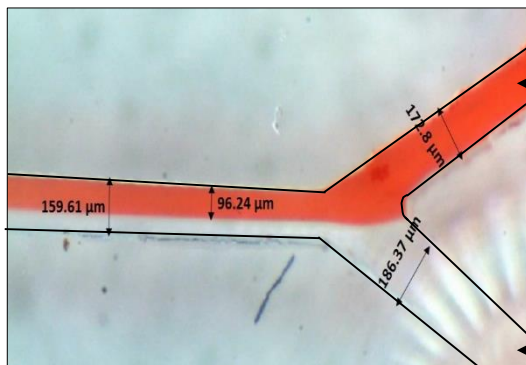
Water with Red ink
at flow rate =
1ml/hr

Water at fixed
flow rate = **1ml/hr**



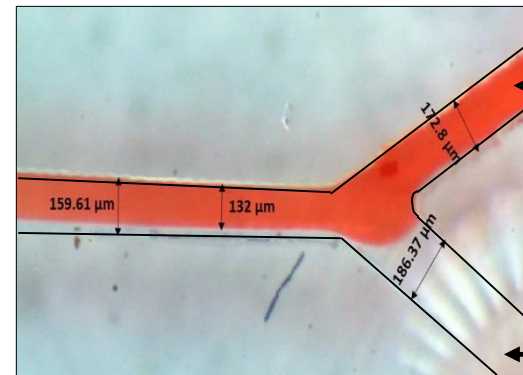
Water with Red
ink at flow rate =
5ml/hr

Water at flow
rate = **1ml/hr**



Water with Red
ink at flow rate =
2ml/hr

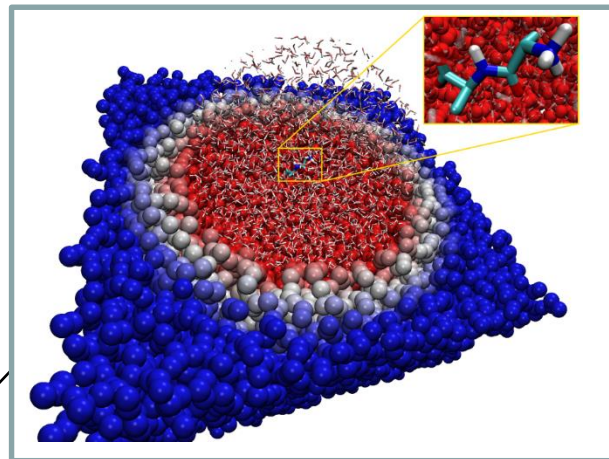
Water at flow
rate = **1ml/hr**



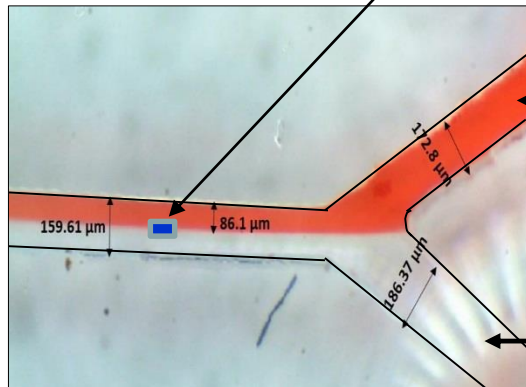
Water with Red
ink at flow rate =
10 ml/hr

Water at flow rate
= **1ml/hr**

Effect of flow rates on protein diffusion and conformational changes

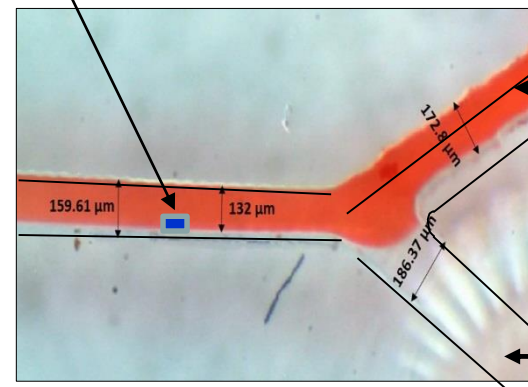


Challenge:
bridging of the time scales?



Water with Red ink
at flow rate =
1 ml/hr

Water at fixed
flow rate = **1 ml/hr**



Water with Red ink
at flow rate =
10 ml/hr

Water at flow rate
= **1 ml/hr**

Multiscale Computing, e.g.

TARDIS=Time Asynchronous Relative Dimension In Space



“Time And Relative Dimensions In Space”
Doctor Who

Journal of Computational Physics 258 (2014) 137–164



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Time asynchronous relative dimension in space method for multi-scale problems in fluid dynamics

A.P. Markesteijn, S.A. Karabasov *



Time Asynchronous Relative Dimension in Space (TARDIS)

- Simple advection equation: $\frac{\partial \rho}{\partial t} + c \frac{\partial \rho}{\partial x} = 0$

$$\bar{x} - x_0 = \alpha (x - x_0) \quad \bar{t} - t_0(t) = \alpha t$$

- Transformation:

$$\alpha = \left(\frac{L_s}{l_s} \right) = \left(\frac{T_s}{t_s} \right)$$

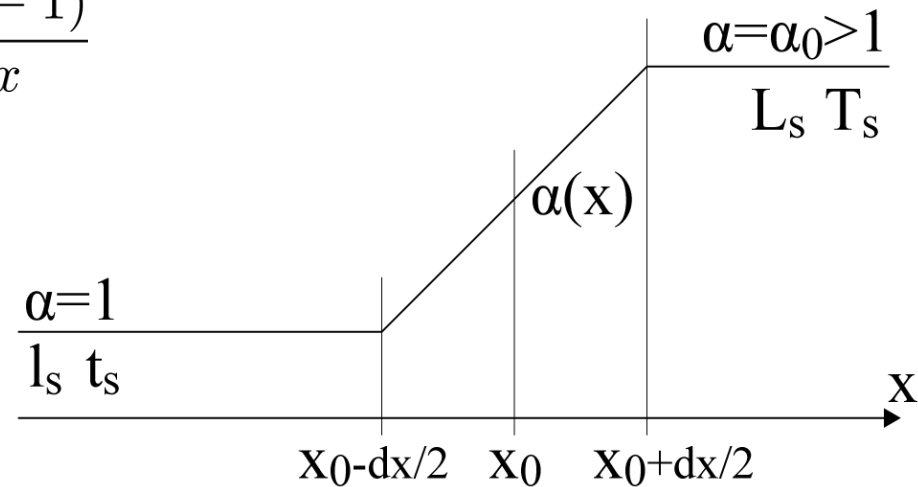
- Introduce time-delay: $\frac{d}{dt} t_0(t) = (x - x_0) \frac{d}{dx} \alpha(x)$

- Where: $\frac{d}{dx} \alpha(x) = \frac{(\alpha_0 - 1)}{dx}$

- Final transformations:

$$\frac{\partial}{\partial x} = \frac{\partial}{\partial \bar{x}} \left(\alpha + (x - x_0) \frac{d}{dx} \alpha \right)$$

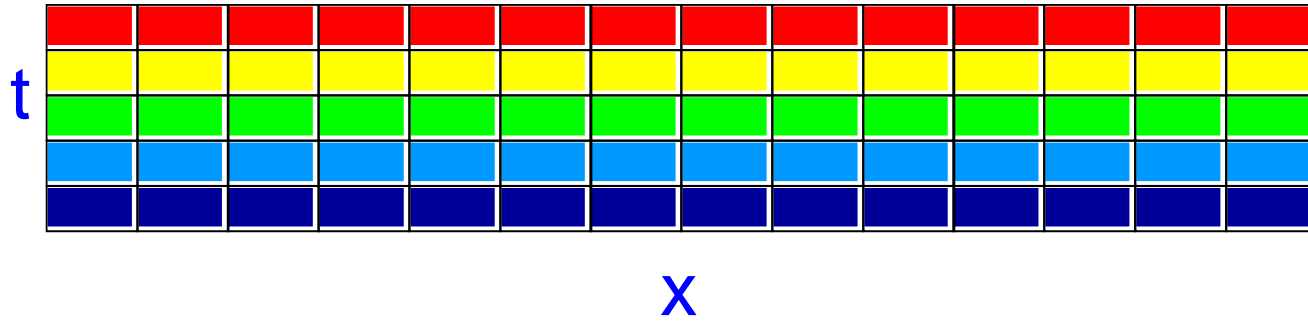
$$\frac{\partial}{\partial t} = \frac{\partial}{\partial \bar{t}} \left(\alpha + (x - x_0) \frac{d}{dx} \alpha \right)$$



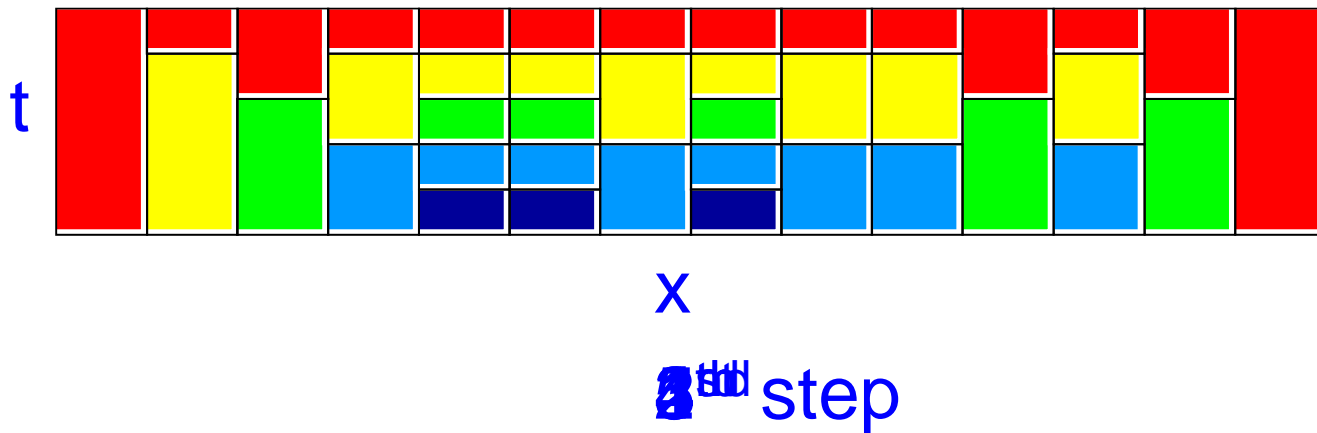
- With compatibility condition:

$$dx \rightarrow 0 : \rho(x, t) \sim \rho(x_0, \bar{t}(x_0) - (\alpha_0 - 1)t) = \rho(x_0, \bar{t}(x_0))$$

Homogeneous time-stepping

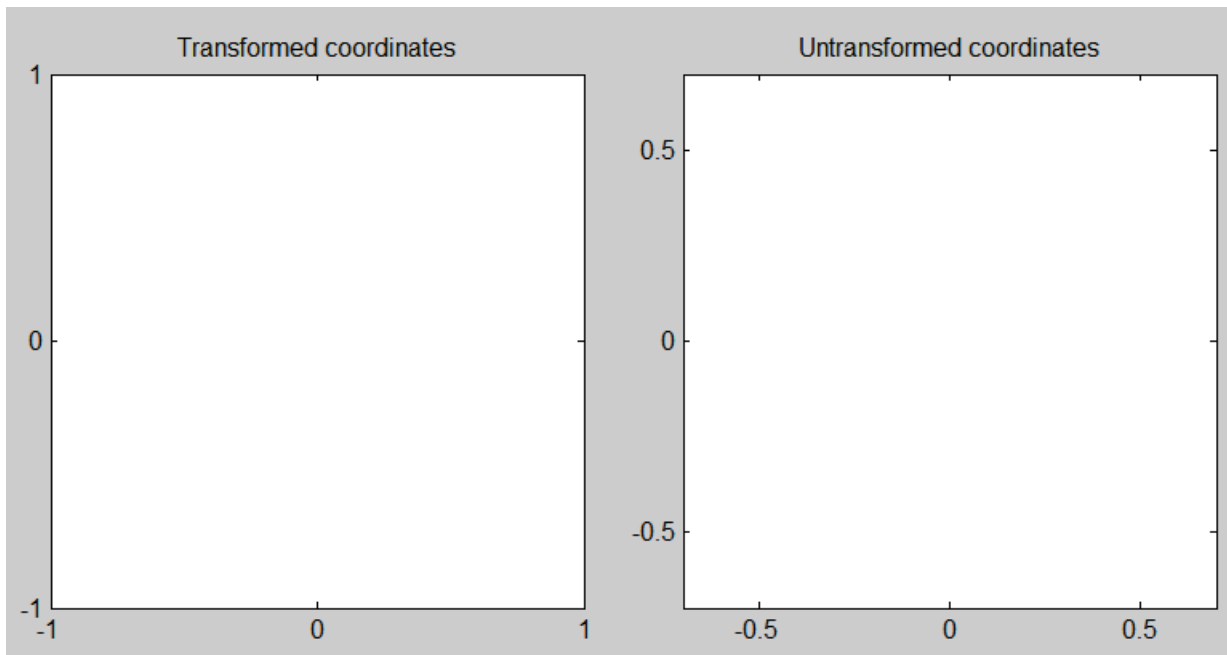
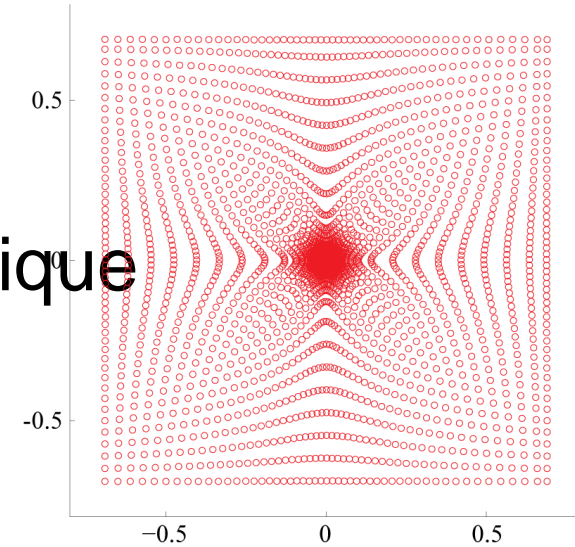


Asynchronous time-stepping



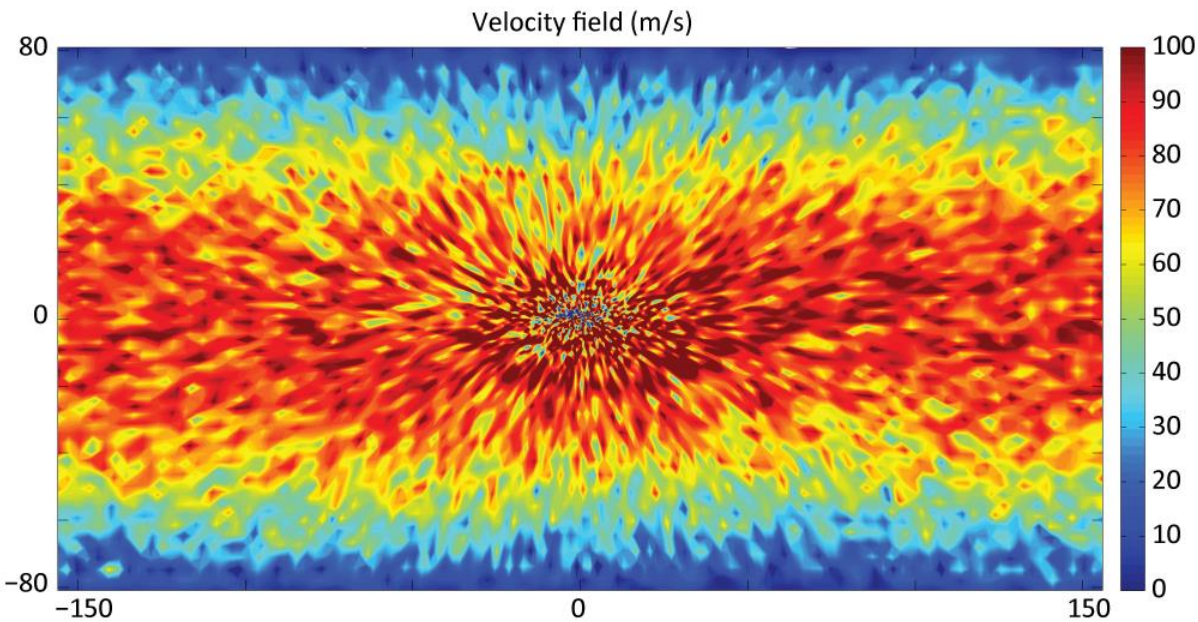
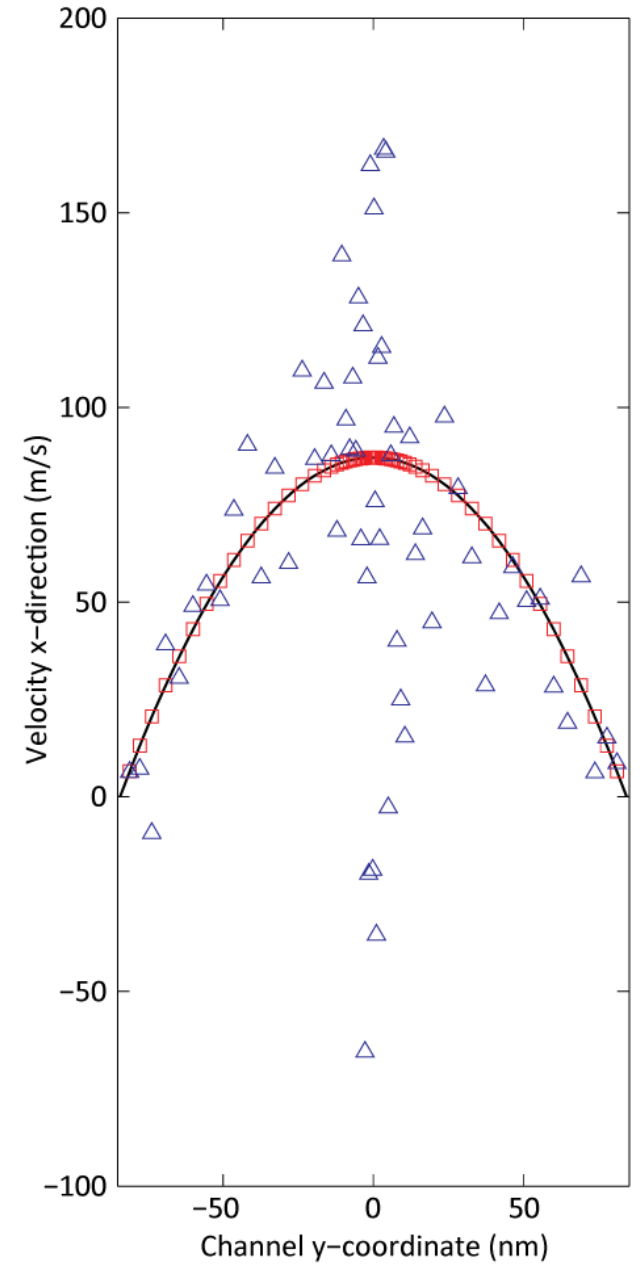
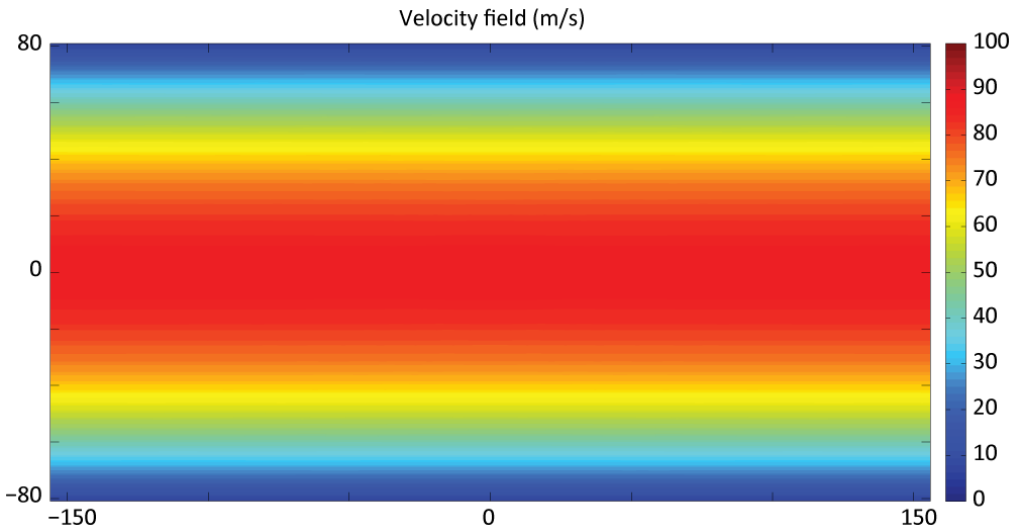
Example: 2D space-time zoom for advection

- Flat boundary equivalent
 - Adjusting aspect ratios of cells
 - Accomplished in ray-tracing technique



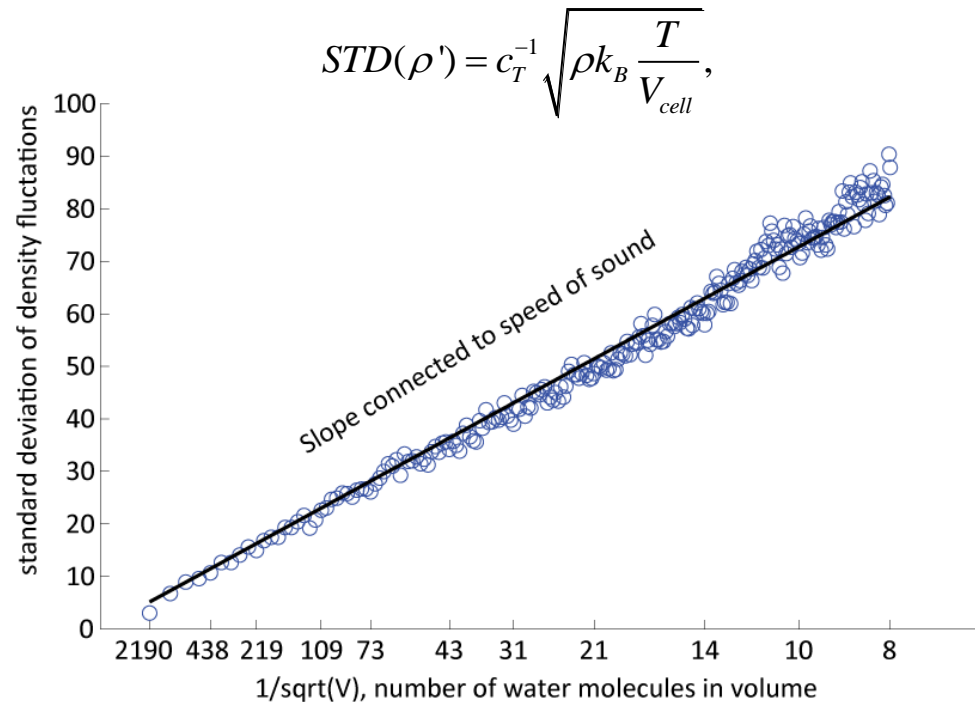
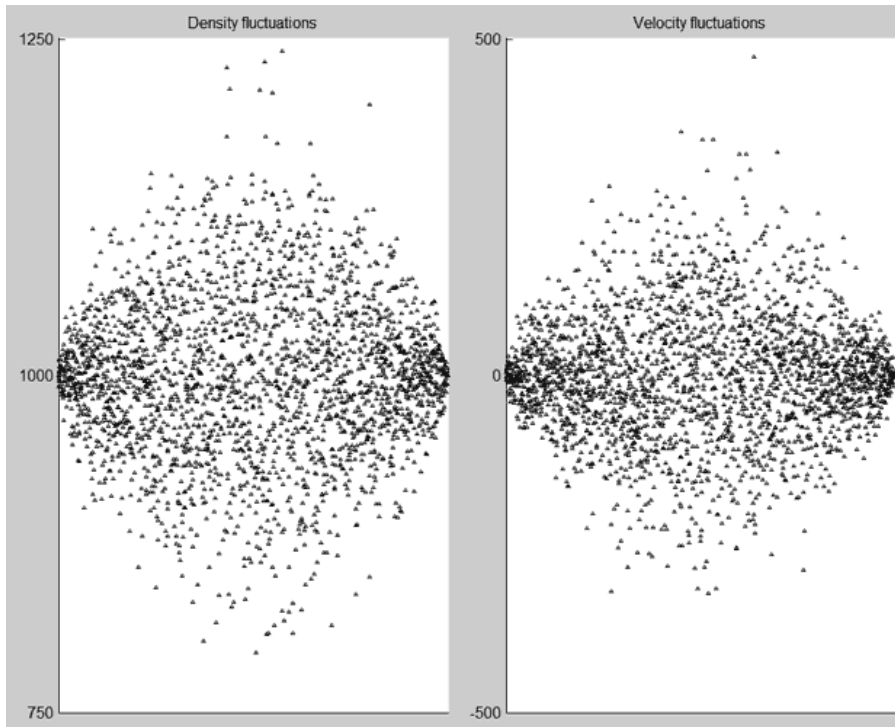
2D channel flow with Fluctuating Hydrodynamics

■ Velocity profile and fluctuations



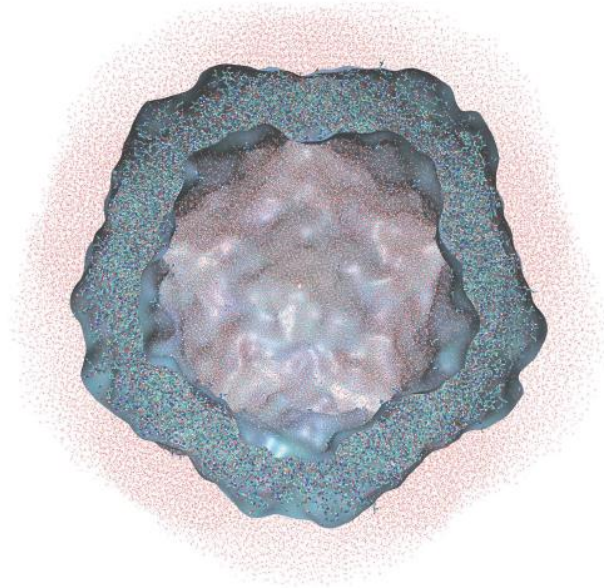
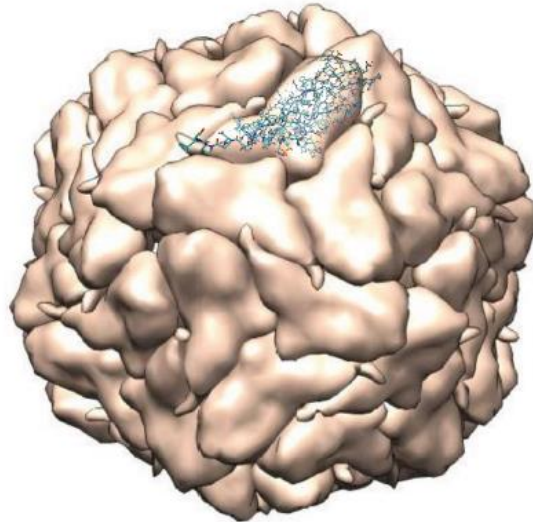
Effect of the scale change on Fluctuating Hydrodynamics

- Density fluctuations and the speed of sound
 - Domain 2560x10, Scale: 1 to 256 to 1 in plateaus
 - Volumes: 2190 → 8 water molecules
 - Speed of sound, relation density, momentum, volume
 - Determined from linear fit: ~1510 m/s (equal to MD results)



Application: simulation of porcine circovirus (PCV2) capsid in water at normal conditions

Capsid diameter of
~18-20 nm (126 180 atoms)
Total box domain size
~ 26.68 nm (~1 180 000 atoms)





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In Press, Corrected Proof — Note to users

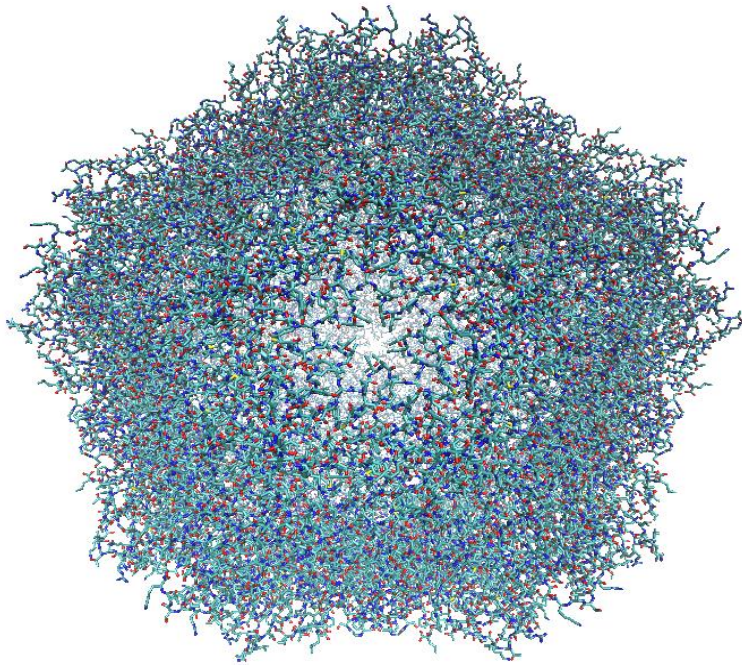


Two-phase flow analogy as an effective boundary condition for modelling liquids at atomistic resolution

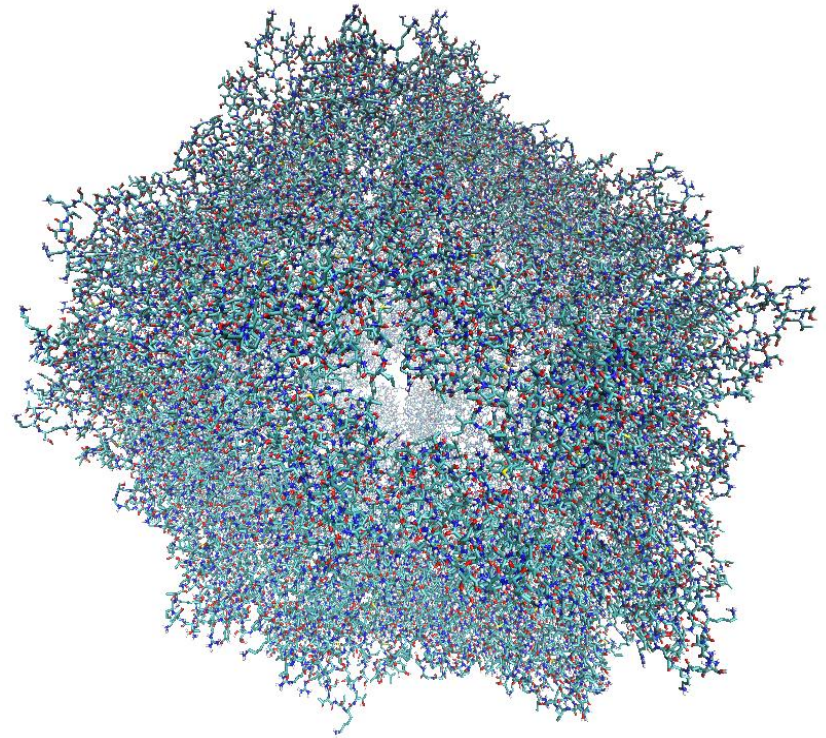
Ivan Korotkin^{a, e}, Dmitry Nerukh^b, Elvira Tarasova^c, Vladimir Farafonov^d, Sergey Karabasov^e.  

[+](#) Show more

Pure MD simulation in a periodic water cube

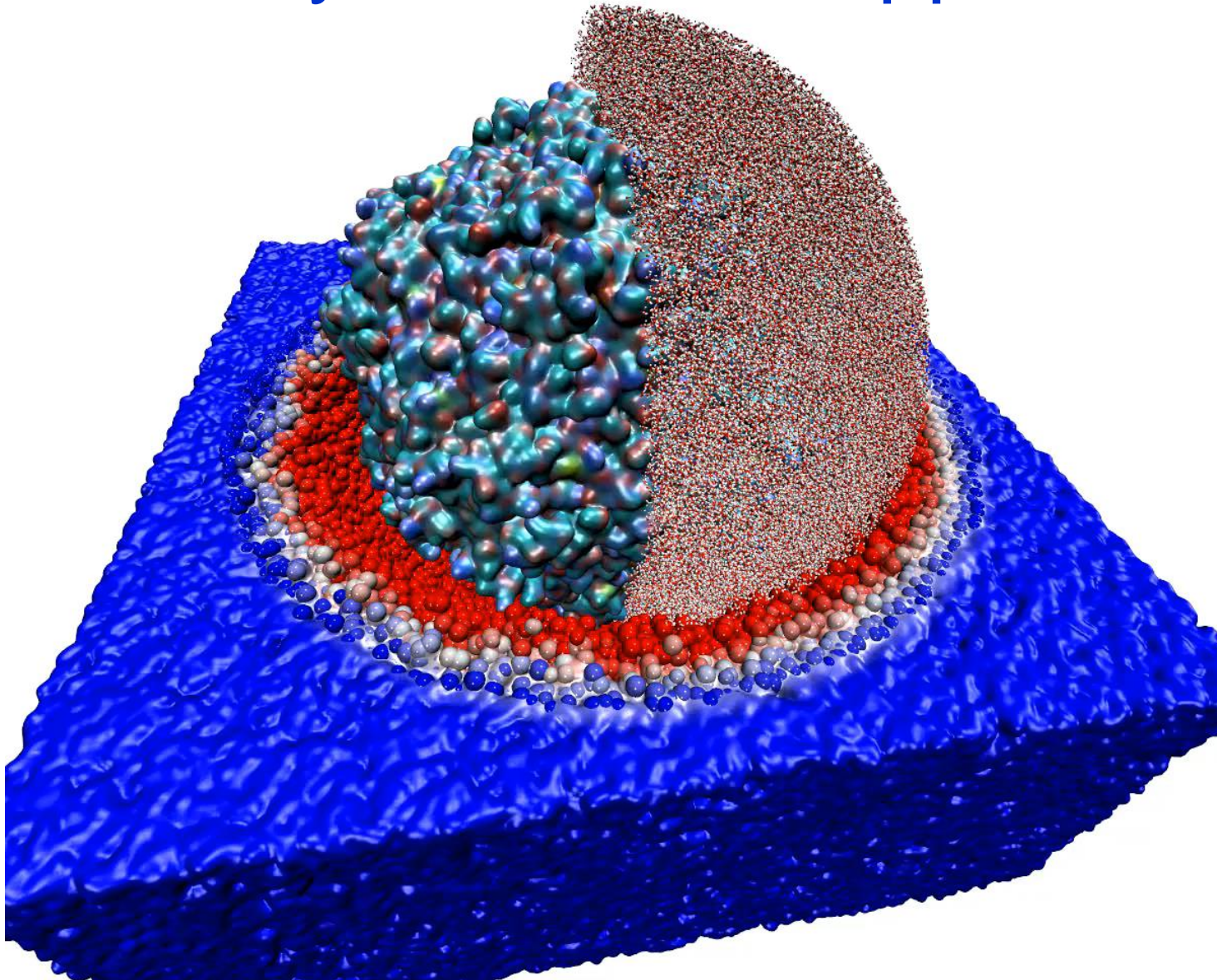


PCV2 capsid initial configuration from the experimental X-ray data



PCV2 capsid becomes unstable in pure MD

Hybrid MD/FH approach

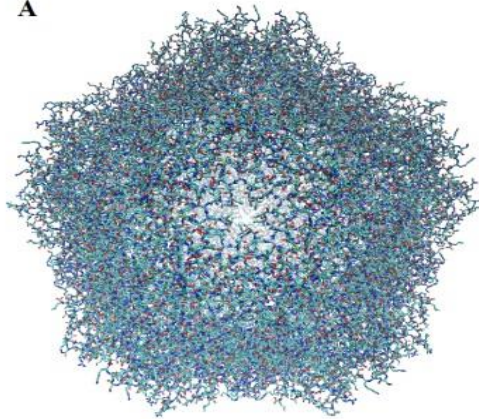


Hybrid MD/FH simulation in a periodic water cube

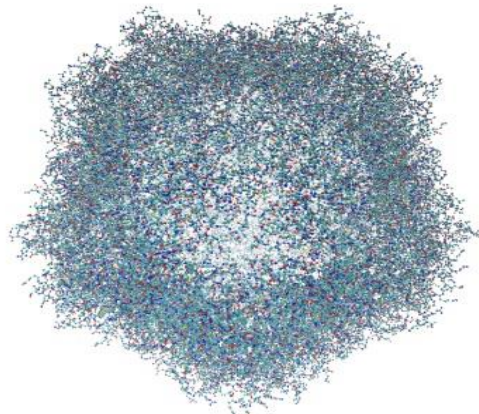
Initial configuration

simulated structure from hybrid MD/FH

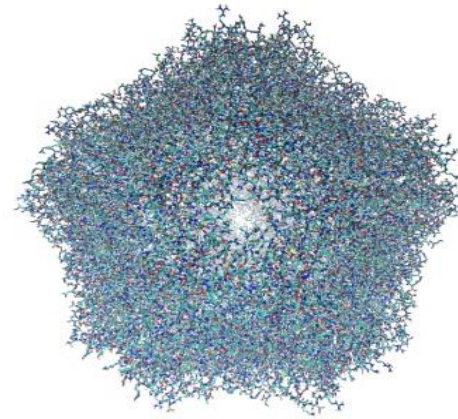
A



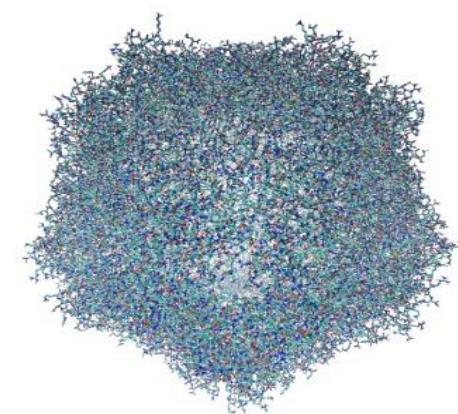
90°



B

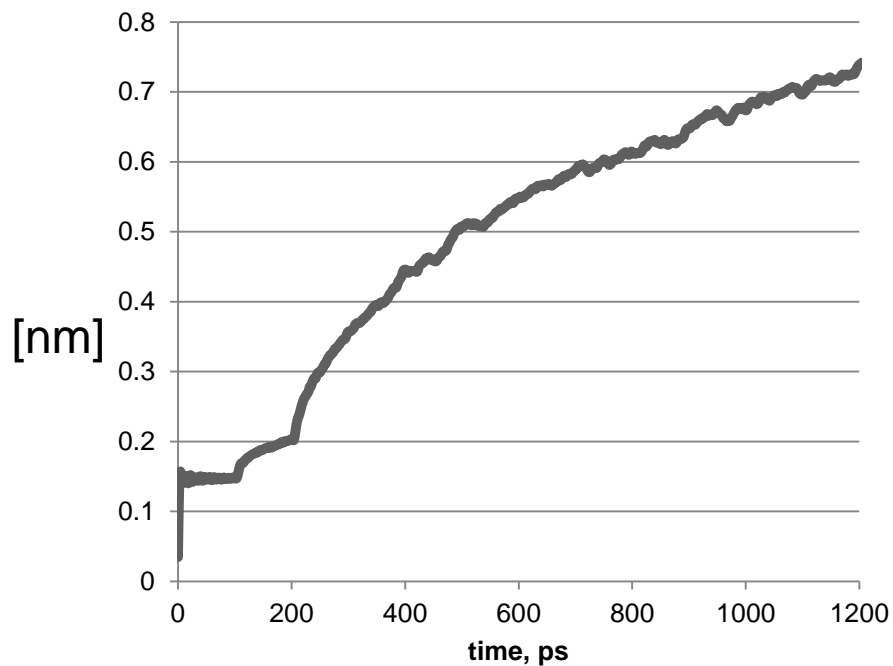


90°

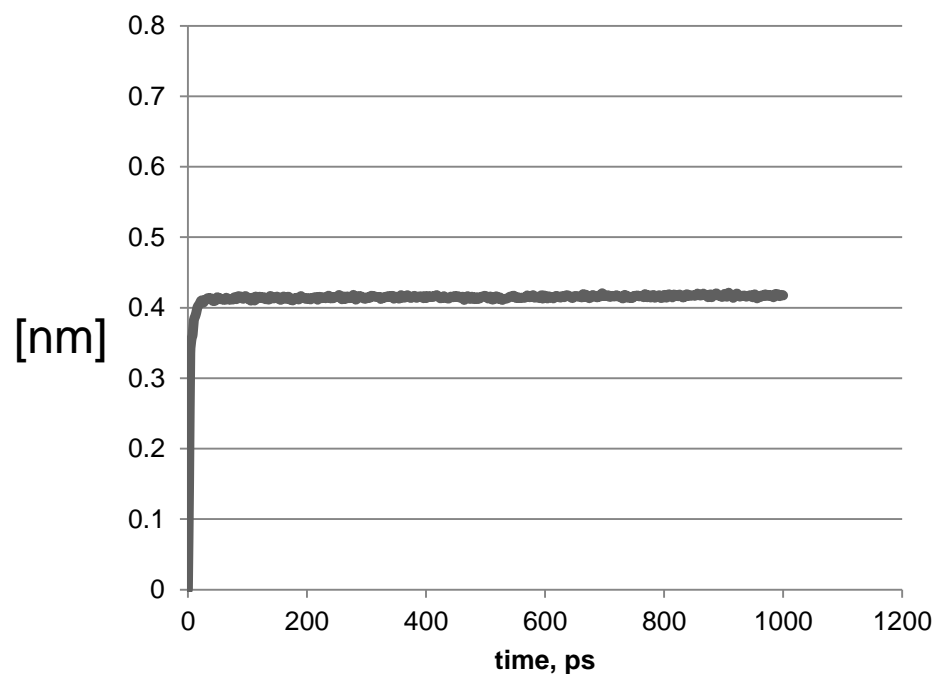


Root-mean-square displacement (RMSD) of the backbone atoms

Pure MD



Hybrid MD/FH



Reference value of RMSD of large molecular systems for stability ~ 0.3nm

Open questions

- A fully coupled MD/FH approach in 3D
- Simulation of large biochemical systems (e.g. more than 1 virus or virus peptide interaction, long simulation times)
- Incorporation of 2d objects (e.g. membranes)
- Adding electrostatic charges to the macroscopic part of the hybrid model
- Generalisations for liquid/solid systems